Casimir Friction

Aleksandr Volokitin

Research Center Jülich and Samara State Technical University



Fluctuation-Dissipation Theorem



$$\mathbf{\Gamma} = (k_B T)^{-1} \operatorname{Re} \int_0^\infty dt \left\langle \hat{\mathbf{F}}(t) \hat{\mathbf{F}}(0) \right\rangle$$

- The detection of single spins by MRFM for: (a) 3D atomic imaging (b) quantum computation
- Measurements of gravitation force at short length scale
- Measurements of Casimir forces

Rytov's Theory



$$\left\langle j_i^f(\mathbf{r}) j_k^{f*}(\mathbf{r}') \right\rangle_{\omega} = \frac{\hbar}{(2\pi)^2} \left(\frac{1}{2} + n(\omega) \right) \omega^2 \mathrm{Im}\varepsilon_{ik}(\mathbf{r}, \mathbf{r}', \omega)$$
$$n(\omega) = [e^{\hbar\omega/k_B T} - 1]^{-1}$$

Origin of Non-Contact Friction



A.I. Volokitin and B.N.J. Persson, *Rev.Mod.Phys.*, **79**, 1291 (2007)

Non-contact Friction



Experiment: Stipe B.C. *et.al PRL*, **87**, 096801 (2001) $F_{friction} = \Gamma v$, $\Gamma \sim 10^{-13} - 10^{-12} \text{kg/s}$ at $d \sim 1 - 100 \text{ nm}$. $\Gamma \approx d^{-n}$ with $n = 1.3 \pm 0.3$ and $n = 0.5 \pm 0.3$, $\Gamma \sim V^2$ Theory: A.I. Volokitin and B.N.J. Persson, *PRB*, **73**, 165423 (2006)

Electronic versus Phononic Friction



Experiment: M. Kisiel et.al Nat. Materials, 10, 119 (2011)

Electronic versus Phononic Friction



Theory: A.I. Volokitin and B.N.J. Persson, *PRB*, **73**, 165423 (2006)

Casimir Friction



Pendry J.B. *JPCM* **9**, 10301 (1997) A.I. Volokitin and B.N.J. Persson, *JPCM*, **11**, 345 (1999)

Frictional Stress

$$\sigma_{xz} = \int_0^\infty \frac{d\omega}{2\pi} \sum_{i=(s,p)} \int \frac{d^2q}{(2\pi)^2} \hbar q_x \operatorname{sgn}(\omega') \left(n_2(\omega') - n_1(\omega) \right) \Gamma_i(\omega, \mathbf{q}),$$

$$\Gamma_i(\omega, \mathbf{q}) = \frac{1}{2k_z \mid 1 - e^{2ik_z d} R_{1i}(\omega) R_{2i}(\omega') \mid^2} \times \left[(k_z + k_z^*)(1 - \mid R_{1i}(\omega) \mid^2)(1 - \mid R_{2i}(\omega' \mid^2) + 4(k_z - k_z^*) \operatorname{Im} R_{1i}(\omega) \operatorname{Im} R_{2i}(\omega') e^{-2d\operatorname{Im} k_z} \right]$$

$$\omega' = \omega - q_x v, \ n_i(\omega) = \left[\exp(\hbar \omega / k_B T_i) - 1 \right]^{-1}, \ k_z = \sqrt{(\omega/c)^2 - q^2}$$

A.I. Volokitin and B.N.J. Persson, JPCM, 11, 345 (1999)

Small Velocities

$$v \ll v_T = k_B T d/\hbar, \ d \ll \lambda_T = c\hbar/k_B T,$$

$$\sigma_{xz} = \gamma v,$$

$$\gamma = \frac{\hbar^2}{8\pi^3 k_B T} \int_0^\infty \frac{d\omega}{\sinh(\hbar\omega/k_B T)} \sum_{i=p,s} \int d^2 q q_x^2 e^{-2qd}$$

$$\times \frac{\mathrm{Im}R_{1i}(\omega)\mathrm{Im}R_{2i}(\omega)}{|1 - e^{-2qd}R_{1i}(\omega)R_{2i}(\omega)|^2},$$

$$\Gamma = \int dS \gamma(z(x, y))$$

Adsorbate enhancement of Casimir friction



A.I. Volokitin and B.N.J. Persson, PRB, 73, 165423 (2006)

Radiative Heat Transfer.



$$T_1 = 273 \text{ K}, T_2 = 0 \text{ K}$$

Volokitin A.I. and Persson B.N.J. PRB, 69, 045417 (2004)

Two ways to study Casimir friction



Left:Upper block is sliding relative to block at the bottom Right: The current is induced in the upper block .

Frictional Drag in 2D-systems



Theory - Coulomb Drag. M. B. Pogrebenskii Sov.Phys.Second.,11 (1977) 372, P. J. Price Physica *B*+*C*,**117** (1983) 750 Experiment - Quantum wells T. J. Gramila et.al PRL,66 (1991) 1216, U. Sivan et.al PRL,68 (1992) 1196 Experiment - Graphene Sheets S. Kim et.al PRB,83 (2011) 161401, R.V. Gorbachev et.al Nat. Phys., 8 (2012) 896 Theory - Casimir Friction. A.I. Volokitin and B.N.J. Persson, J.Phys.:Condens.Matter, 13, 859 (2001); ibid, EPL 103 24002 (2013)

Quantum Friction



$q_x v = \omega_1 + \omega_2$ Pendry J.B. *JPCM* **9**, 10301 (1997) Volokitin A.I. and Persson B.N.J. *JPCM*, **11**, 345 (1999) Volokitin A.I. and Persson B.N.J. *PRB*, **78**, 155437 (2008)

Anomalous Doppler Effect



Normal Doppler effect $\omega' = \omega - q_x v > 0$ Anomalous Doppler effect $\omega' = \omega - q_x v < 0$ Thermal fluctuations dominate at $v < v_T = k_B T d/\hbar$ Quantum fluctuations dominate at $v > v_T = k_B T d/\hbar$

Classical Vavilov-Cherenkov Radiation



Cherenkov P.A. Dokl. Akad. Nauk SSSR, **2**, 451 (1934) Resonant condition: $q_x v = cq/n > v_0 = cq_x/n$ Threshold velocity: v > c/n

Quantum Vavilov-Cherenkov Radiation



Frank M.I. J.Phys.USSR, 7, 49 (1943) Doppler effect (a): $\omega_{ph} = \omega_0 - g_x v > 0$ Doppler effect (b): $\omega_0 - g_x v < 0$; $\omega_{ph} = q_x v - \omega_0$

The Landau criterion for the critical velocity of a superfluid



$$E = \frac{Mv^2}{2} + \varepsilon(p) - pv$$

$$\varepsilon(p) - pv < 0$$

$$v > v_c = min\left(\frac{\varepsilon(p)}{p}\right)$$

Landau L.D., *Zh. Eksp. Teor. Fiz.* 11, 592 (1941)

Radiation at shearing two transparent plates



Pendry J.B. *JMO* **45**, 2389 (1998). Magreby F.M., Golestian R., and Kardar M., *PRA* **88**, 042509 (2013). Volokitin A.I. and Persson B.N.J. *PRB* **93**, 035407 (2016).

Non-Relativistic Theory



Resonant condition: $\omega_{ph} = q_x v - cq/n = cq/n, v > v_0 = 2c/n$ $F_{1x}^s = \frac{\hbar v_0}{\pi^3 d^4} F_{1x}^s$ [normalized]

$$R_s = \frac{\sqrt{(\omega/c)^2 - q^2} - \sqrt{(n\omega/c)^2 - q^2}}{\sqrt{(\omega/c)^2 - q^2} + \sqrt{(n\omega/c)^2 - q^2}}$$

Magreby F.M., Golestian R., and Kardar M., *PRA* 88, 042509 (2013). Volokitin A.I. and Persson B.N.J. *PRB* 93, 035407 (2016).

Relativistic Theory



$$\omega_{ph} = q_x v - cq' / \gamma n = cq/n, v > v_0 = 2cn/(n^2 + 1)$$

Volokitin A.I. and Persson B.N.J. *PRB* **78**, 155437 (2008) Volokitin A.I. and Persson B.N.J. *PRB* **93**, 035407 (2016).

Radiation From Moving Neutral Particle



$$\omega_{ph} = q_x v - \omega_0 / \gamma = cq/n, v > v_0 = c/n$$

Pieplow G. and Henkel C., *JPCM* **27**, 035407 (2015). Volokitin A.I. and Persson B.N.J. arxiv.org/abs/1512.04366 (2016).

Quantum Friction for Particle



Polar Dielectric SiO₂



Resonant condition: $\omega_{ph} = q_x v - \omega_0 = \omega_0$ Threshold velocity: $v > 2\omega_0/q_{xmax} \sim 2\omega_0 d \sim 2 \cdot 10^5$ m/s

$\textbf{Graphene on SiO}_2$



Threshold velocity: $v > v_F + \omega_0/q_{xmax} \approx v_F \sim 10^6$ m/s Resonant condition: $\omega_{ph} = q_x v - v_F q = \omega_0$ Experiment:Freitag M.,Steiner M., Martin Y., Perebeinos V., Chen Z., Tsang J.C., and Avouris P., Nano Lett. 9, 1883 (2009). Theory:Volokitin A.I. and Persson B.N.J. PRL 106 094502 (2011)

Current density-electric field dependence in graphene on SiO₂



$$v_{sat} \sim v_F \sim 10^{\circ} m/s$$

 $J_{sat} = en_s v_{sat} \sim 1 \, mA/\mu m$

Frictional Drag between Graphene Sheets



d=1 nm d=10 nm At I $v \ll v_F$ induced electric field $E = \rho_D J = \mu^{-1} v$.

$$\rho_D = \frac{\Gamma}{(ne)^2} = \frac{h}{e^2} \frac{\pi\zeta(3)}{32} \left(\frac{k_B T}{\epsilon_F}\right)^2 \frac{1}{(k_F d)^2} \frac{1}{(k_T F d)^2},$$
$$F_{x0} = \frac{hv}{d^4} \frac{15\zeta(5)}{128\pi^2} \left(\frac{v}{v_F}\right)^2 \frac{1}{(k_T F d)^2}.$$

Volokitin A.I. and Persson B.N.J. EPL 103 24002 (2013)

Acoustic Phonon Emission



$$\sigma(\mathbf{x},t) = K[u_0(\mathbf{x} - \mathbf{v}t, t) - u_1(\mathbf{x}, t)],$$

Persson B.N.J., Volokitin A.I. and Ueba H., *JPCM* **23** 045009 (2011)

Acoustic Phonon Emission



Summary

- Casimir friction can be studied using modern experimental setups for observation of frictional drag in graphene systems..
- The challenging problem is to study frictional drag for graphene for strong electric field (large drift velocity) when Casimir friction is dominated by quantum fluctuations (quantum friction).
- The challenging problem for experimentalists is to develop setup for mechanical detection of Casimir friction using AFM.
- Quantum friction is strongly enhanced above the threshold velocity.
- The threshold velocity is determined by a type of excitations which are responsible for quantum friction.