

Near-field thermal effects at mesoscopic scale

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Outline:

sub-Kelvin to Kelvin temperature range
near-field heat transfer

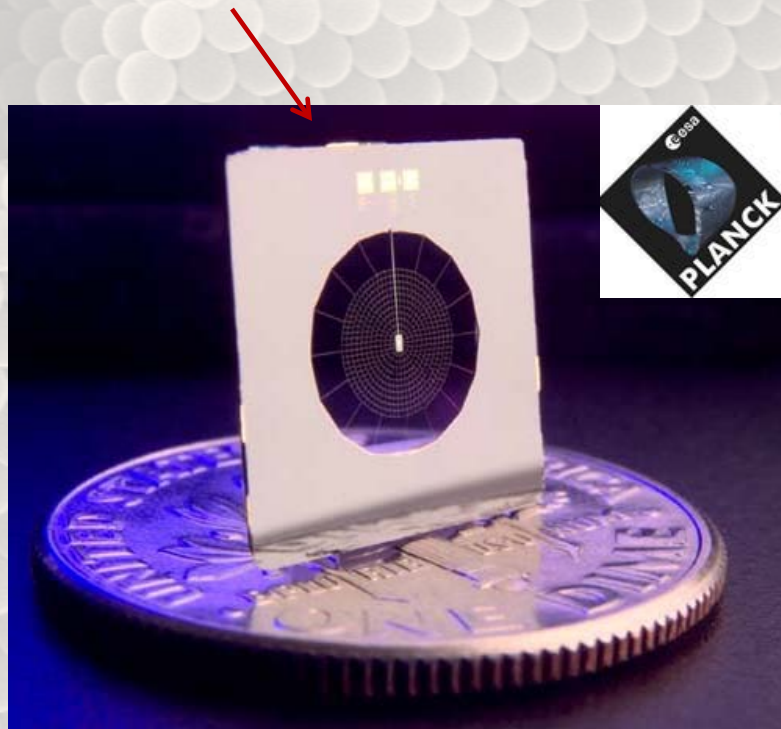
- Simple theory
- Experimental design
- Preliminary data

Why low temperatures?

- Can access extreme near field conditions more easily, as dominant thermal wavelengths are 2-3 orders of magnitude larger
- Possible relevance in low-temperature detector applications, where *typical* operating temperatures are ~ **0.1 K**

Examples of ultra-sensitive devices at low temperatures:

Spider-web bolometers, force/mass NEMS detectors, transition edge sensors



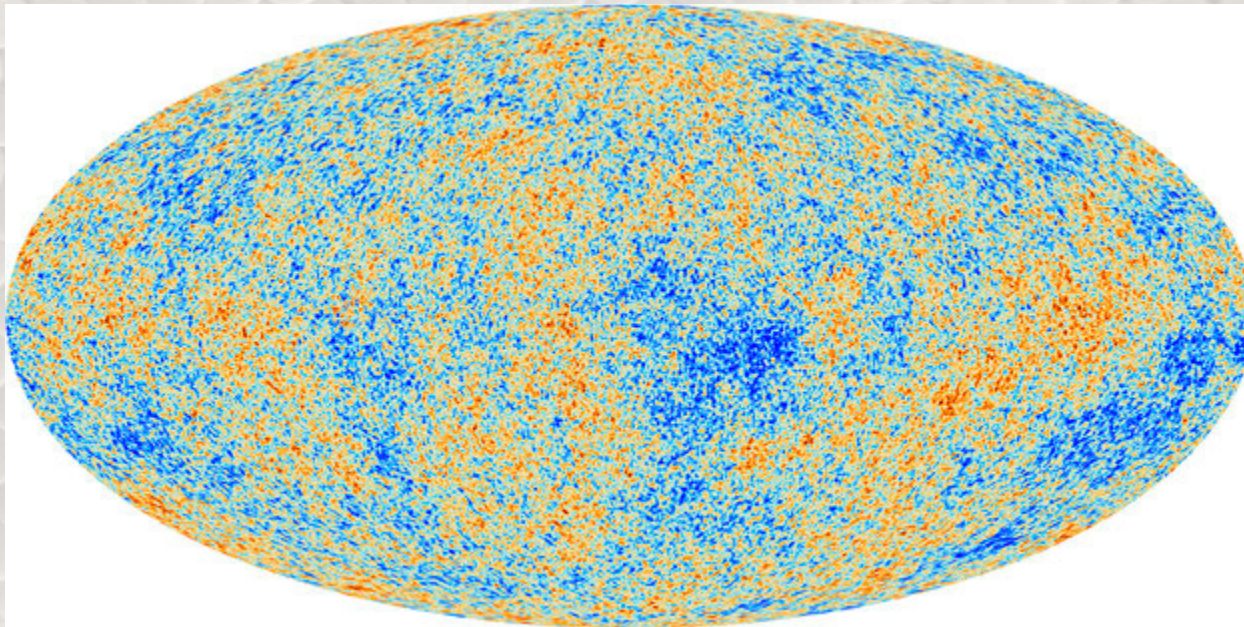
JPL built 0.1 K spider-web bolometer in Planck

G phonon ~ 100 pW/K , P ~ 2.5 pW

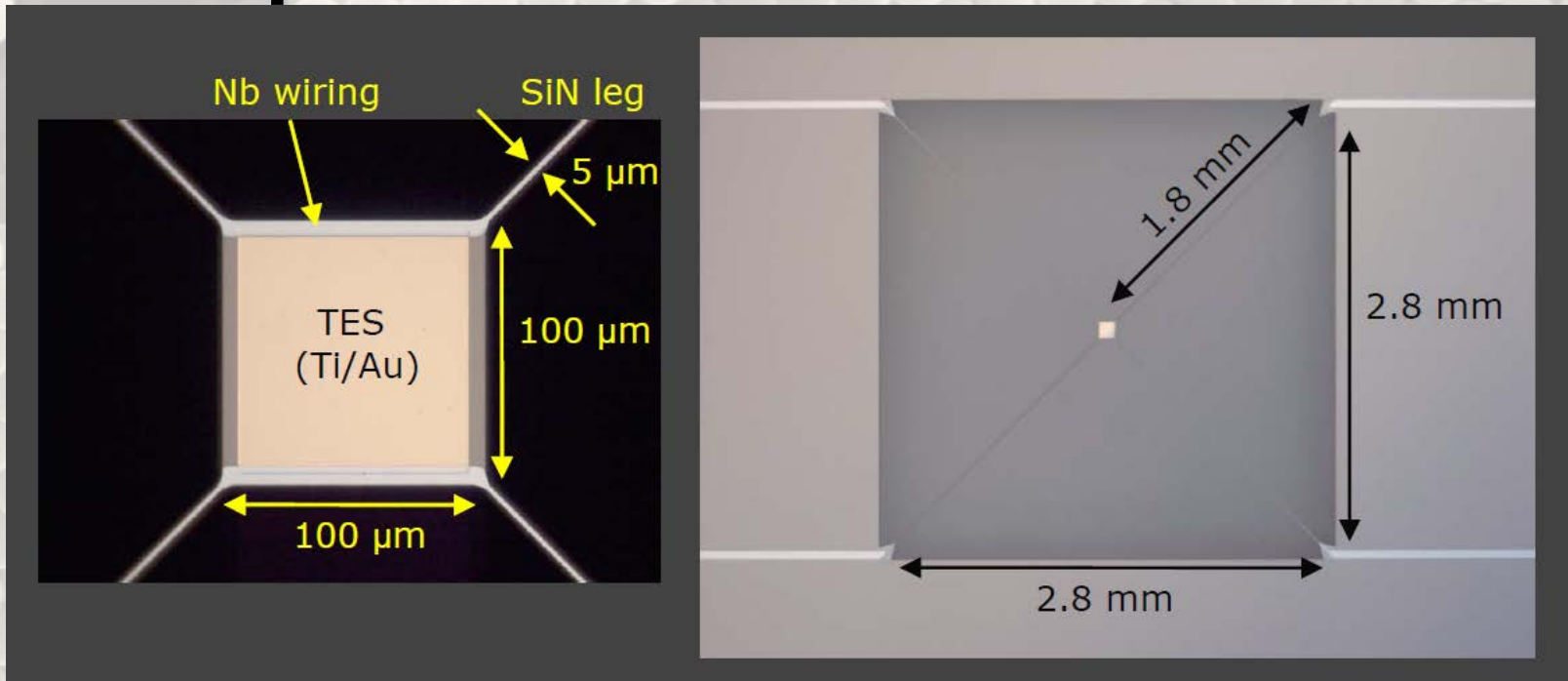


0.1 K TES X-ray detector array and SQUID readout, NIST+JYU

Without low-T detectors, no high-res CMB data !
Space is too HOT for the detectors (dilution refrigerator in space!)



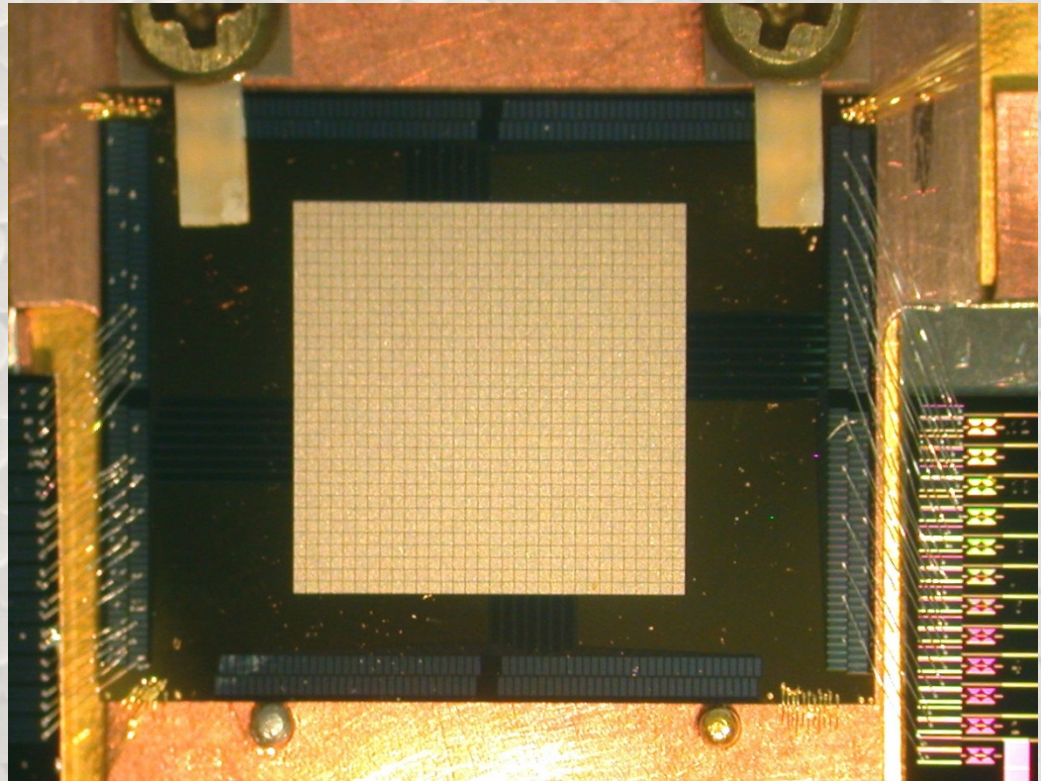
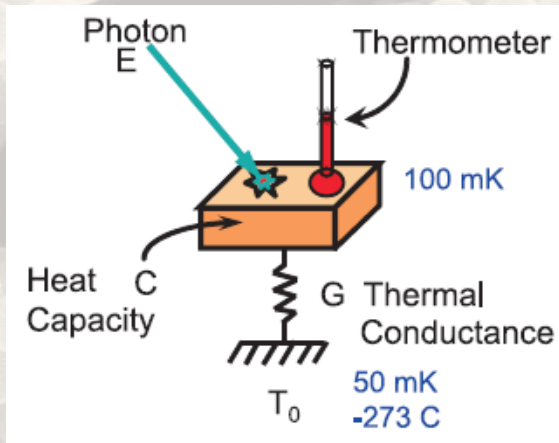
Example *real* devices in more detail



SRON FIR bolometer for future astronomy mission (SPICA) suspended SiN beams requirement

$G \sim 0.1 \text{ pW/K}$, $P \sim 3 \text{ fW}$

Example *real* devices in more detail 2:



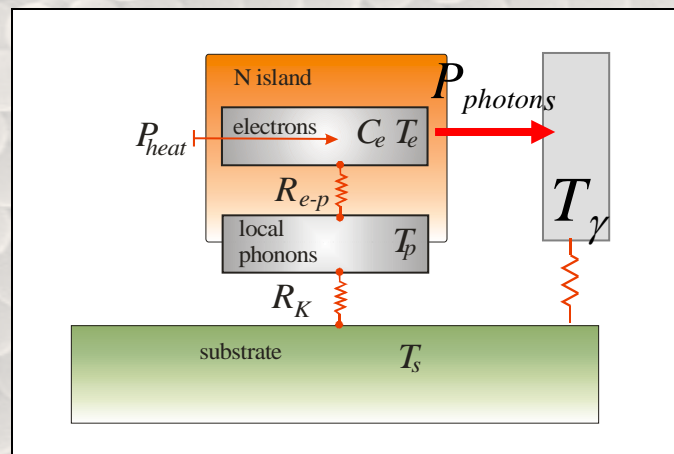
$G \sim 300 \text{ pW/K}$
 $P \sim 50\text{-}100 \text{ pW/pixel}$

Thermal cross-talk??

NASA Goddard $32 \times 32 = 1024$ pixel superconducting X-ray transition edge sensor array
Finnish 256 pixel array in progress (Jyväskylä+VTT)

=> Need to understand and control thermal conductance in nanoscale

For bolometers, if thermal conductance is low, small heat loads lead to large temperature increase => more sensitivity



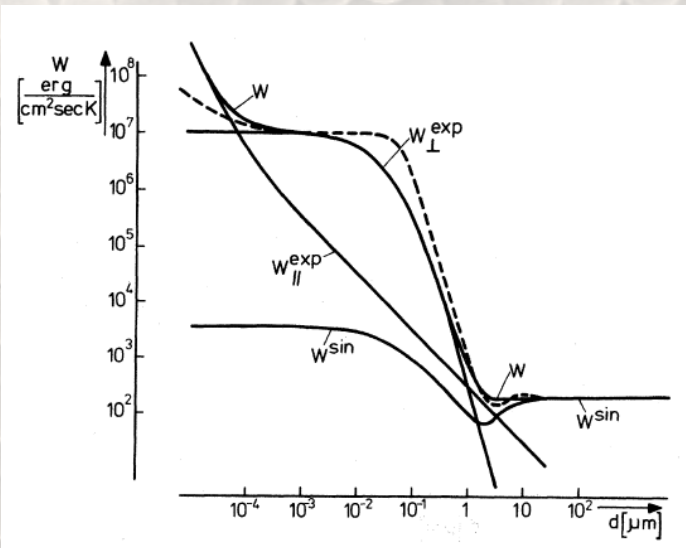
Thermal model for samples

- Electron-phonon interactions
- Phonon heat conductance
- Photon heat conductance

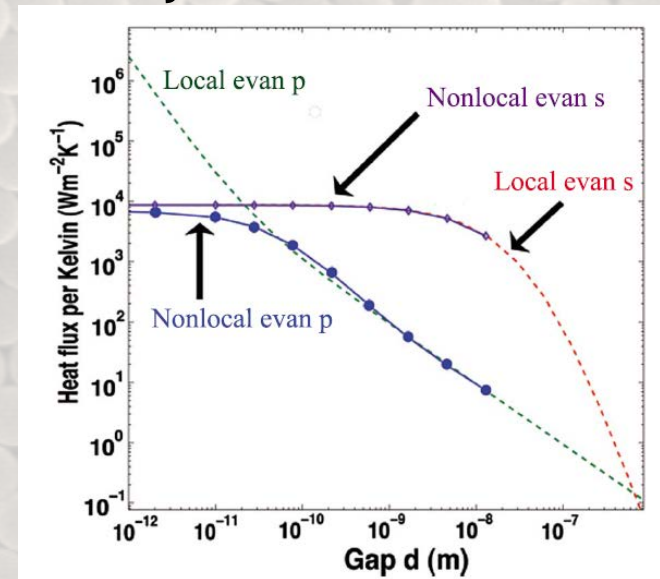
Low thermal conductance and cooling **increases bolometer performance**
($NEP \sim G^{1/2}T$)

How does near-field thermal transport depend on temperature in an ideal case (Drude metals, parallel planar surfaces)?

Most publications discuss only RT results

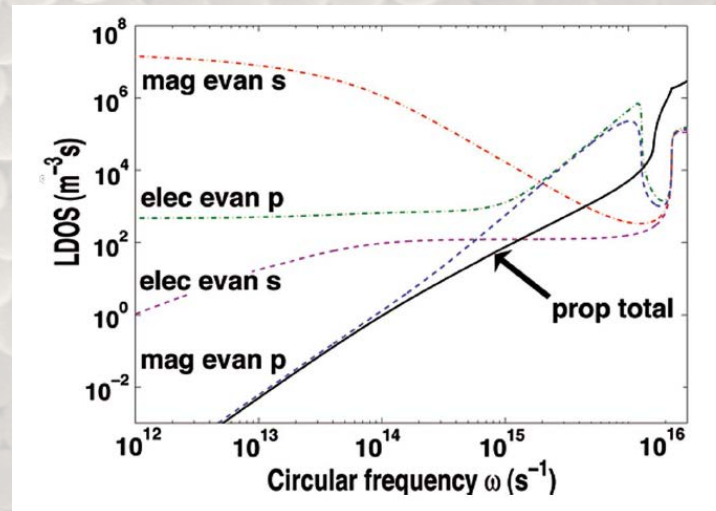


Polder and Van Hove, PRB 1971



Chapuis, Volz, Henkel, Joulain, Greffet, PRB 2008

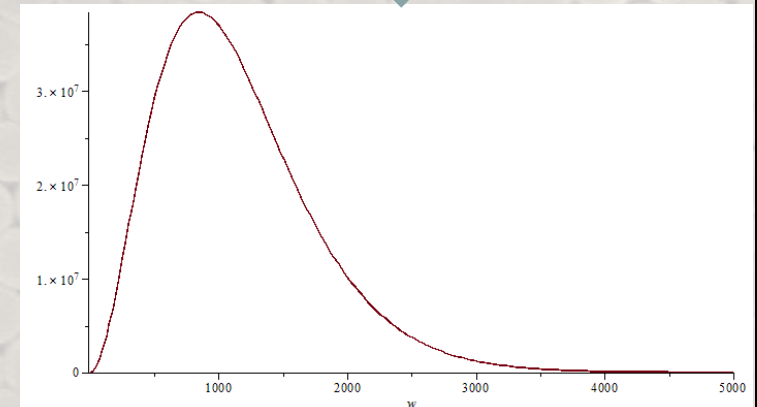
- For a typical Drude metal (Au,Cu,Al...) **at RT:**
- Thermal near field starts at \sim **few μm**
- Is dominated by s (TE)-polarized evanescent waves, with magnetic fields dominating
- Reaches saturation at distance \sim 1 nm
- Has a maximum power enhancement $\sim 10^5$



Temperature acts as a low-pass filter, cutting off higher frequency components, cut-off frequency moves linearly up with T (Wien's law)

Planck spectrum

$$I_{\omega}^0 = \frac{\omega^2}{4\pi^3 c^2} \frac{\hbar\omega}{(e^{\hbar\omega/k_B T} - 1)}$$

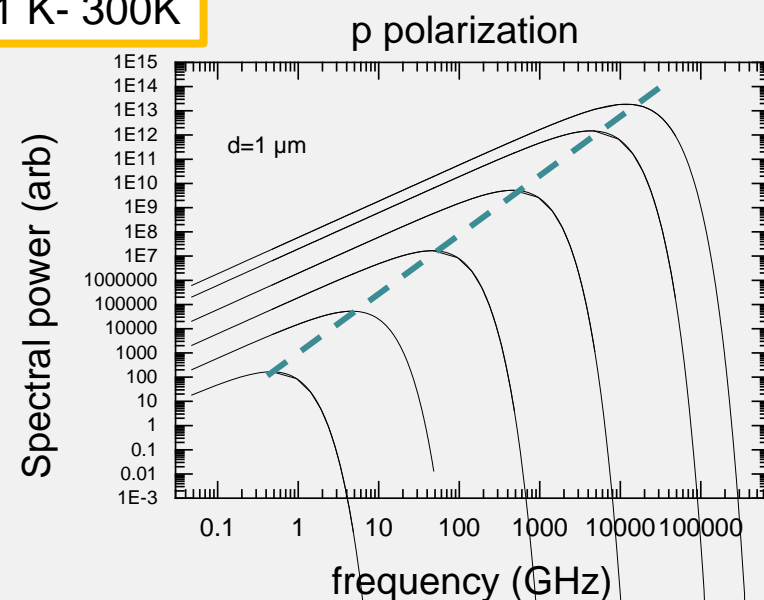
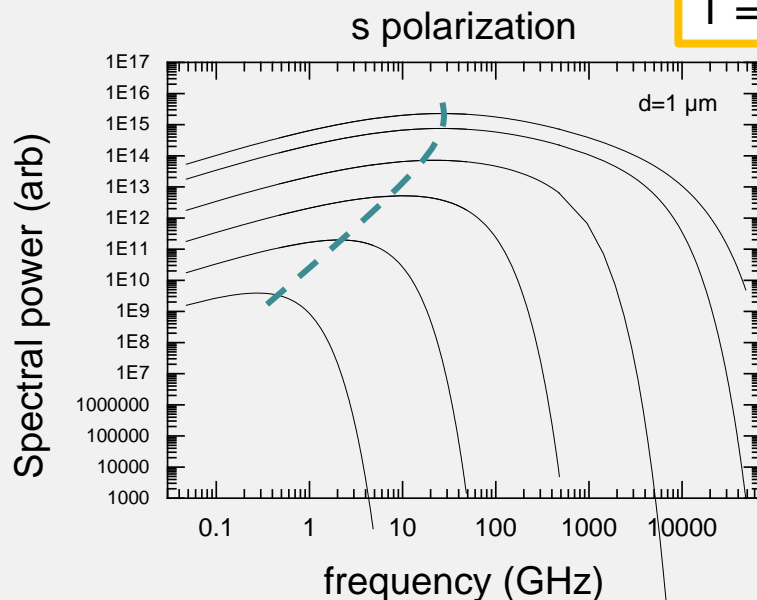


$$\begin{aligned} \phi = & \int_{\omega=0}^{+\infty} d\omega [I_{\omega}^0(T_1) - I_{\omega}^0(T_2)] \\ & \times \sum_{\alpha=s,p} \left[\int_0^{\omega/c} \frac{KdK}{\omega^2/c^2} \frac{(1 - |r_{31}^{\alpha}|^2)(1 - |r_{32}^{\alpha}|^2)}{|1 - r_{31}^{\alpha} r_{32}^{\alpha} e^{2i\gamma_3 d}|^2} \right. \\ & \left. + \int_{\omega/c}^{\infty} \frac{KdK}{\omega^2/c^2} \frac{4 \operatorname{Im}(r_{31}^{\alpha}) \operatorname{Im}(r_{32}^{\alpha}) e^{-2\gamma_3'' d}}{|1 - r_{31}^{\alpha} r_{32}^{\alpha} e^{-2\gamma_3'' d}|^2} \right], \end{aligned}$$

Example calculations for Drude Cu with a measured low-T mean free path:

Power spectra are very different from Planckian for s-polarization (but close for p polarization)

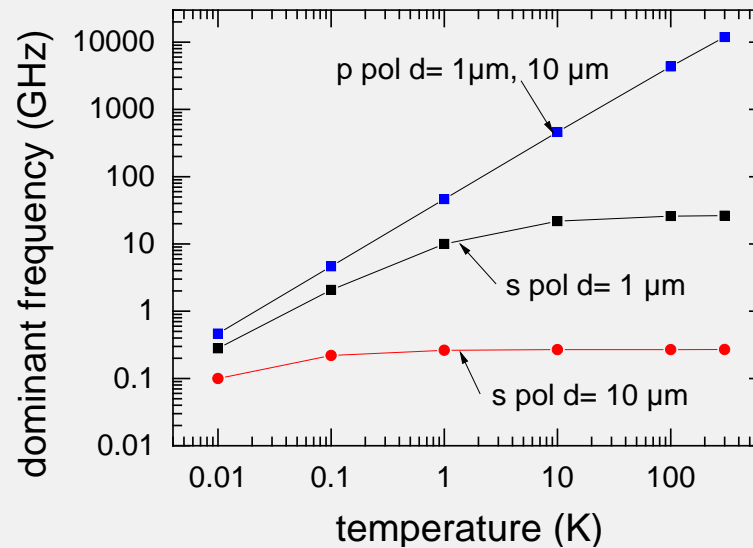
$T = 0.01 \text{ K} - 300 \text{ K}$



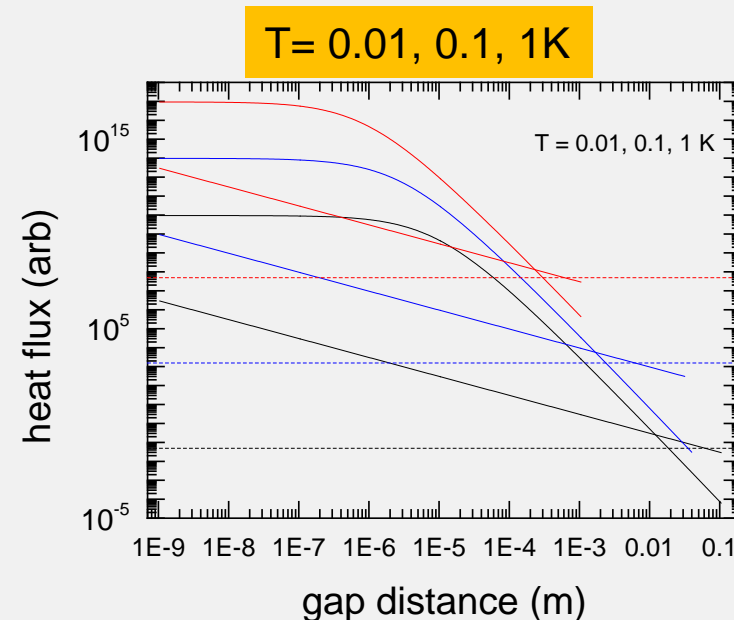
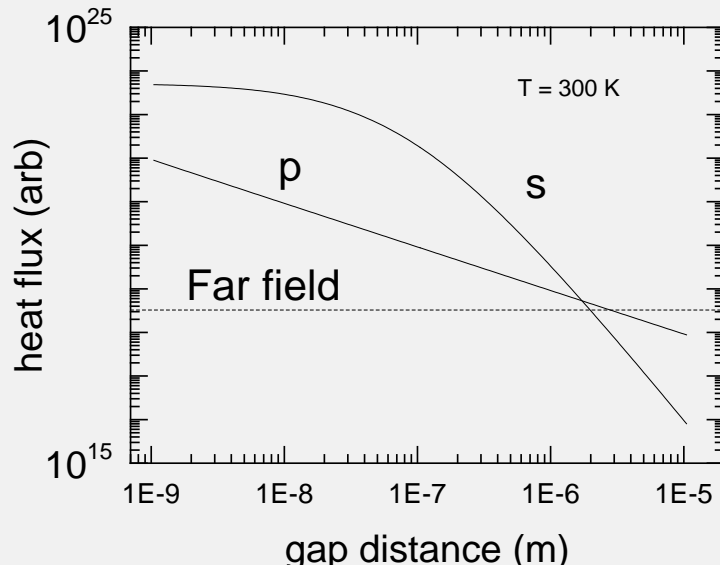
Dominant frequency does not scale linearly with T for s-polarization !
(3 orders of magnitude difference at RT)

Position of dominant frequency for s-polarization also strongly dependent on gap distance

$T = 0.01 \text{ K} - 300 \text{ K}$

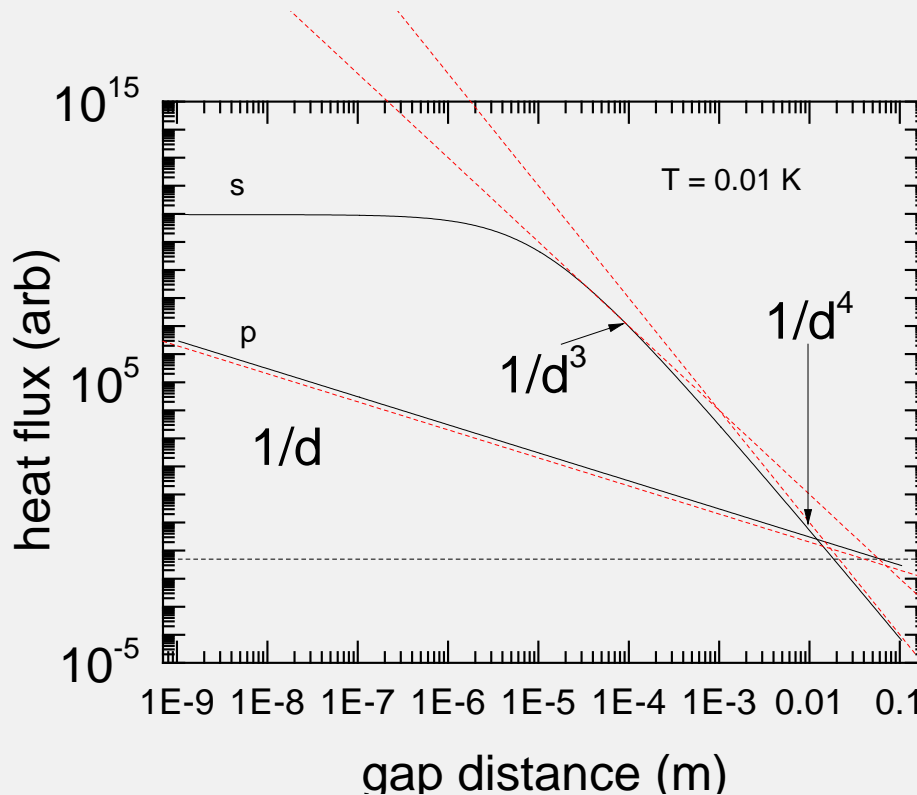


Example calculations for Drude **Cu** with a measured low-T mean free path:

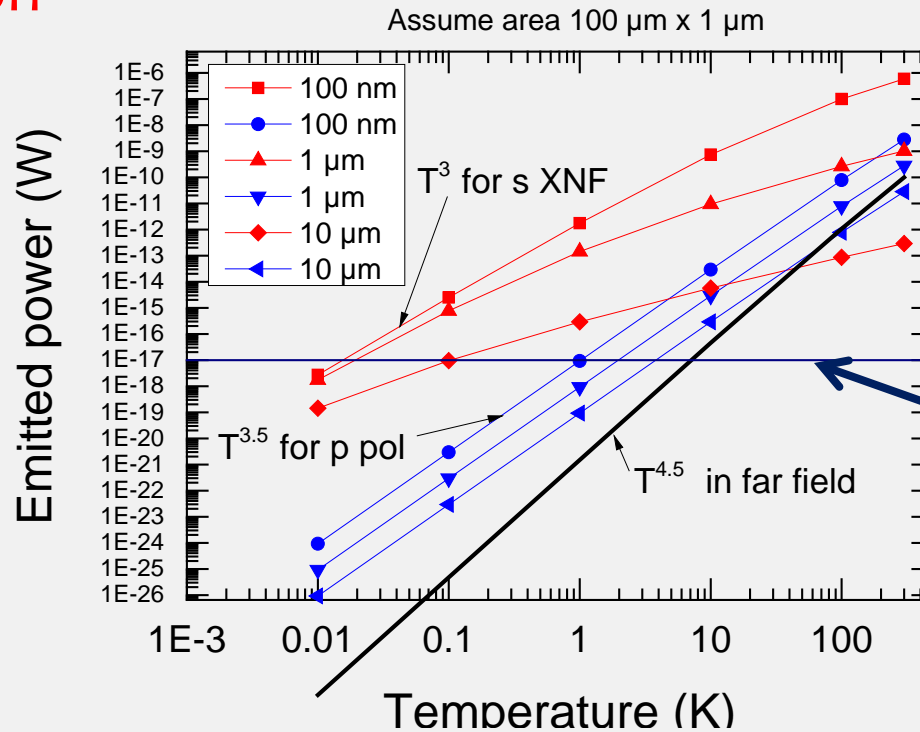


- Thermal near field starts at $\sim 1 \text{ mm} - 10 \text{ cm}$ distance
- Is still dominated by s polarized evanescent waves, but a window of p-dominance appears at low T when power first starts increasing
- Reaches saturation at distance $\sim 1 \mu\text{m}$ (0.01 K) – 100 nm (1K)
- Has a maximum power enhancement $\sim 10^{12}$ (0.01 K) – 10^9 (1 K)

- Distance dependence is not affected by temperature



- Power enhancement is stronger because **near field contribution dies out more slowly with T than the far-field contribution**

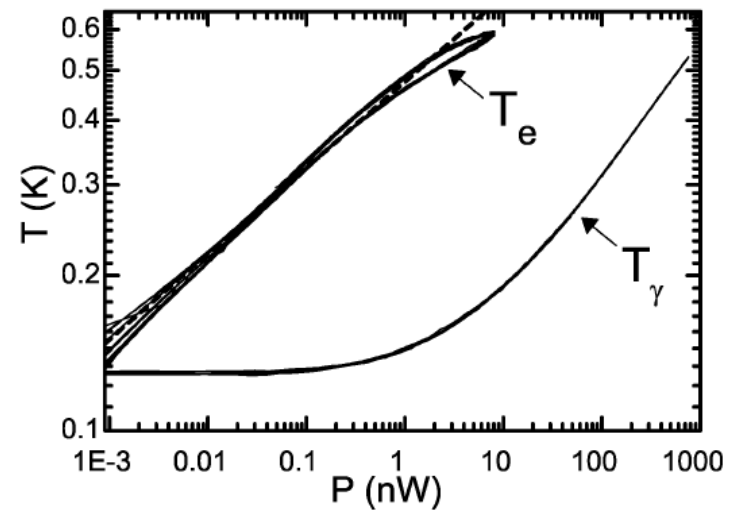
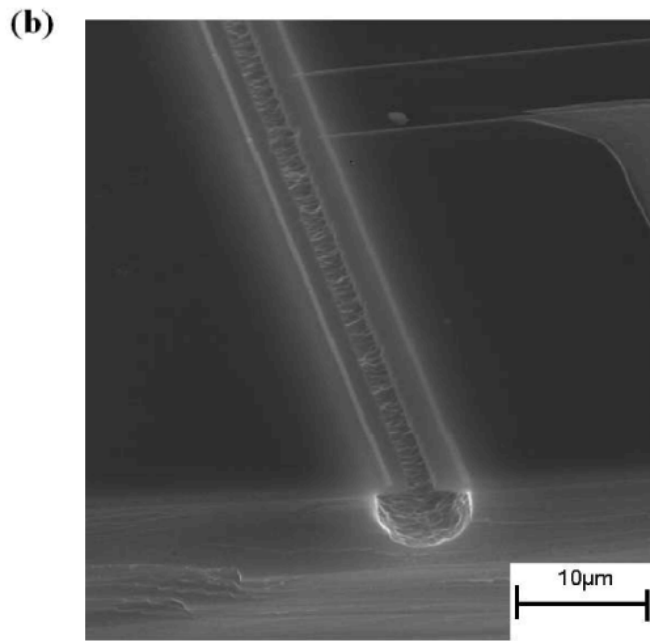


Current measurement limit for our structures

Near field heat transfer measurable at mesoscopic distances (up to $\sim 10 \mu\text{m}$) at cryogenic temperatures !

- sub-Kelvin to Kelvin*** temperature range
near-field heat transfer
- Experimental design
 - Preliminary data

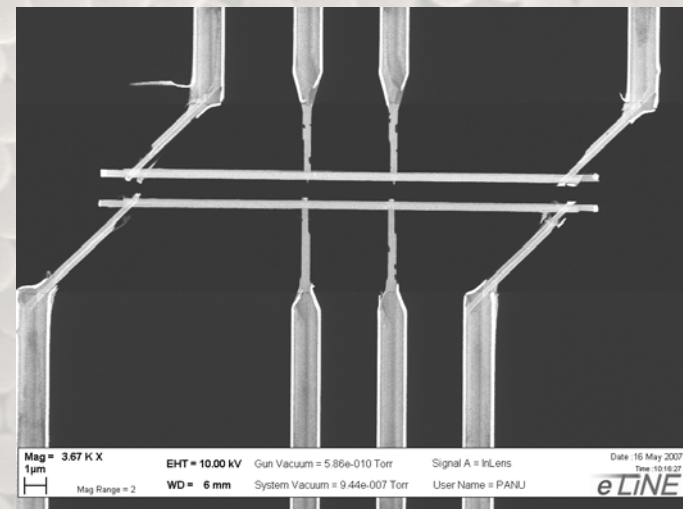
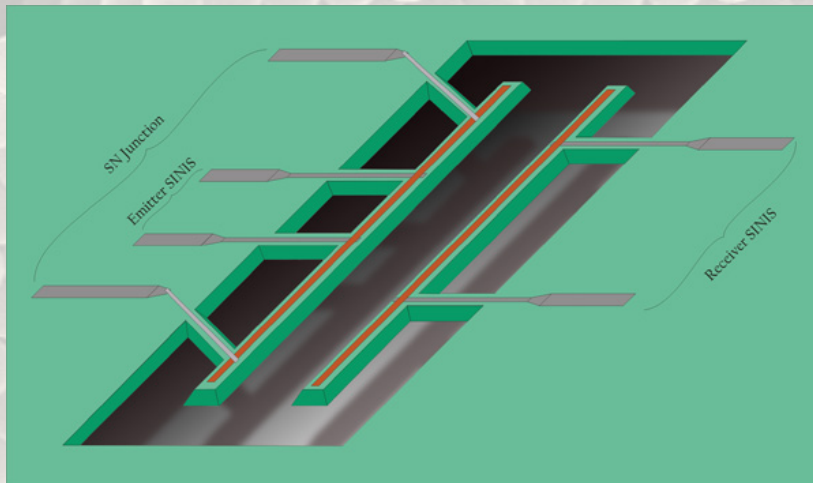
Early idea: etch a trench between two metallic wires



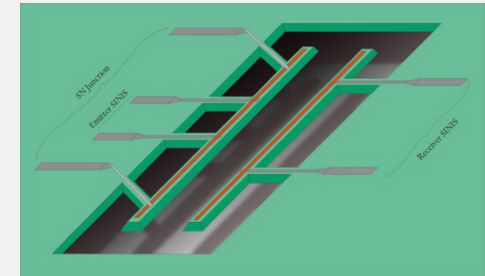
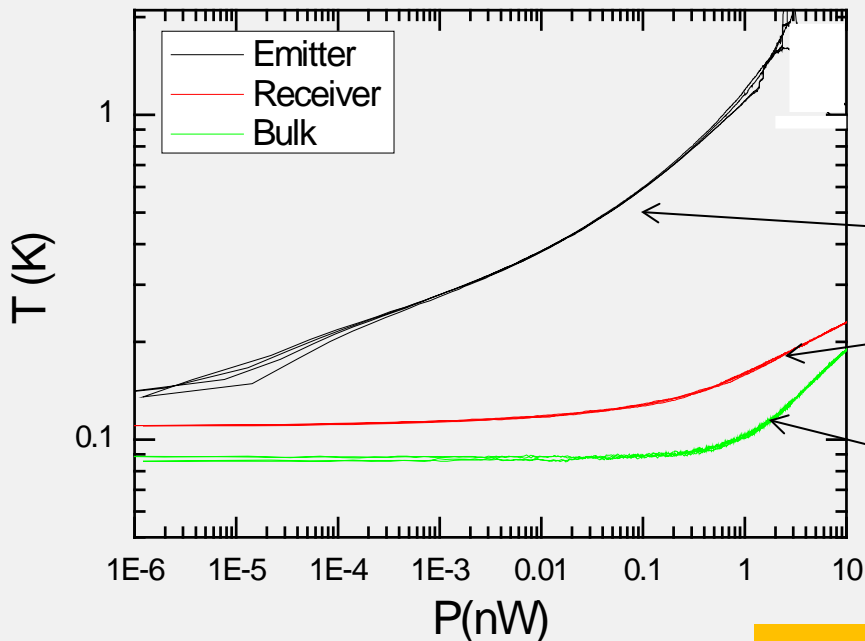
BUT: can we be sure that phonon conduction
Still doesn't dominate ??

P.J. Koppinen, J. T. Karvonen, L. J. Taskinen, I.J. Maasilta, *AIP Conf. Proc.* **850**, 1556 (2006)

=> Move to suspended wire geometry to fully remove substrate between wires



Again clear heating signal in the receiver wire



Heated wire

Clear signal in the receiver wire

Much weaker signal in phonon thermometer

BUT: what is EMITTED NF power ?

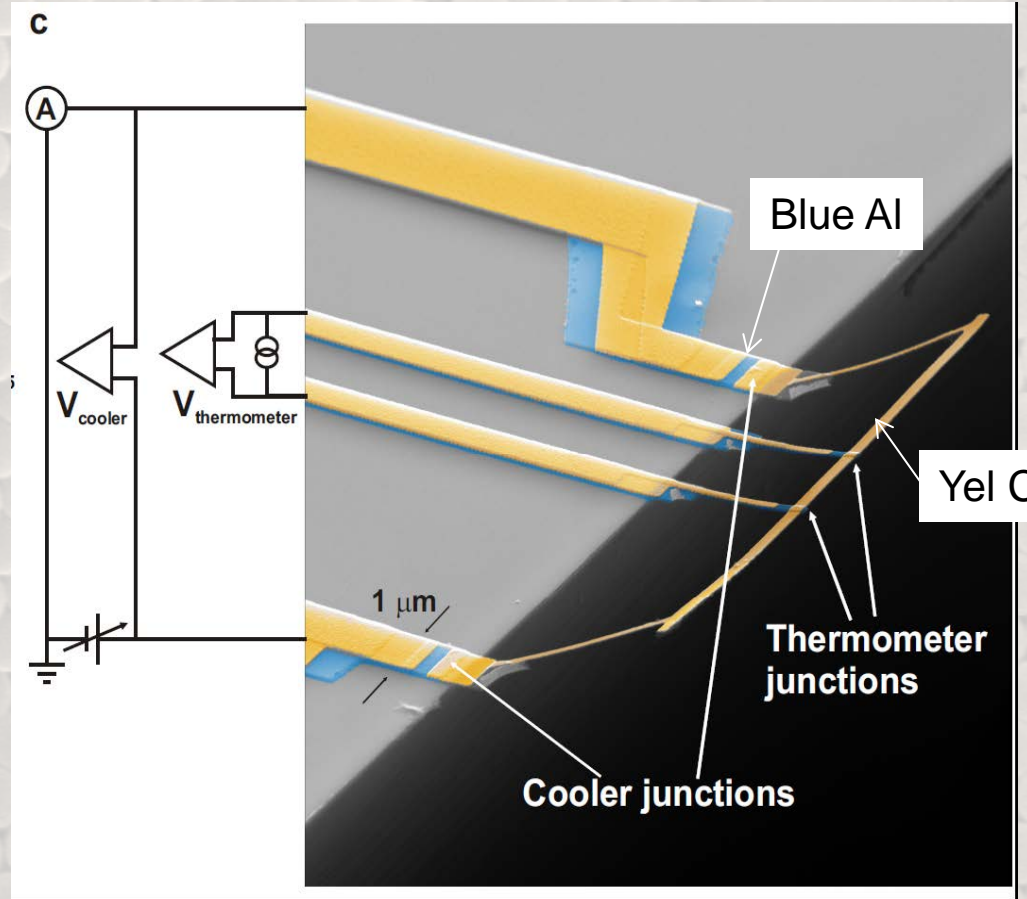
P.J. Koppinen, T. J. Isotalo, I.J. Maasilta, *TRN 07, Les Houches (2007)*

Low G (phonon thermal conductance) due to nanoscale *suspended beams*

- SINIS
(Superconductor-Insulator Normal metal) tunnel junction thermometry $< 1\text{K}$
- SINIS tunnel junction *coolers*

nanowire length 10-20 μm ,
thickness 60 nm and width 150-300 nm
4 supporting bridges length 5 μm ,
thickness 60 nm and width 150 nm

P.J. Koppinen, I.J. Maasilta,
Phys. Rev. Lett. **102**, 165502 (2009)



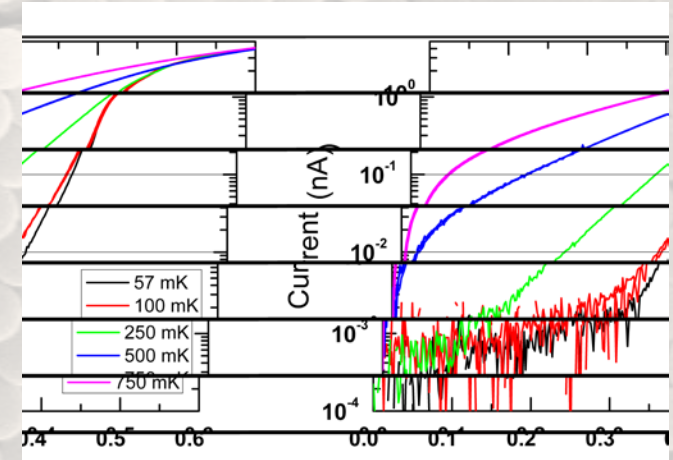
Tunnel junction thermometry

I-V characteristics non-linear with temperature

$$I(V) = \frac{1}{eR_T} \int_{\Delta}^{\infty} \frac{|E|}{\sqrt{E^2 - \Delta^2}} [f(E - eV, T_e) - f(E + eV, T_e)] dE$$

Independent of superconductor temperature

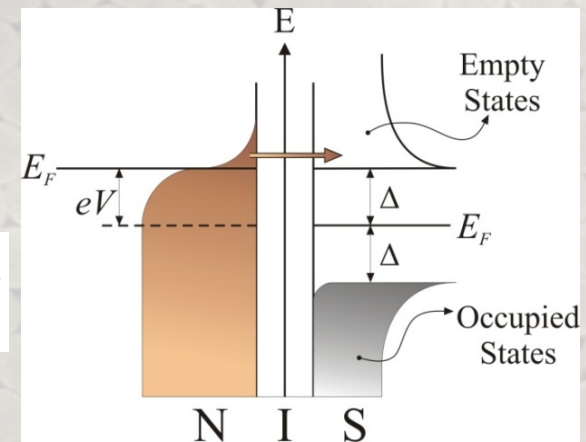
Resolution (~ 0.1 mK at 100 mK)



Tunnel junction cooling

Tunneling of “hot” electrons from Fermi tail
(bias voltage dependent, optimal at $V \sim \Delta$)

$$\dot{Q}_{cool} = \frac{1}{e^2 R_T} \int_{-\infty}^{\infty} (E - eV) g_S(E) [f_N(E - eV, T_N) - f_S(E, T_S)] dE$$

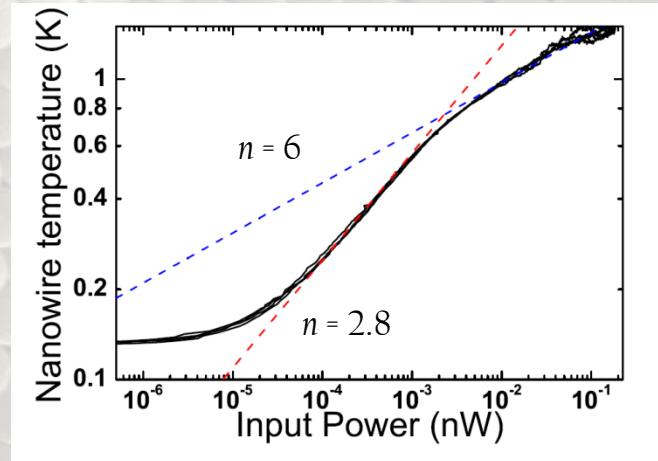


direct heating experiment without coolers

$n=2.8$ consistent with 1D-2D interface scattering [1]

No T-gradients within wire =>

Heat flow dominated by the nanowire-bulk interface



$$G = dP/dT = 0.4 \text{ pW/K at } 0.2 \text{ K (} 0.1 G_Q / \text{channel)}$$

Phonon emission power ~ 50 fW at 0.2 K

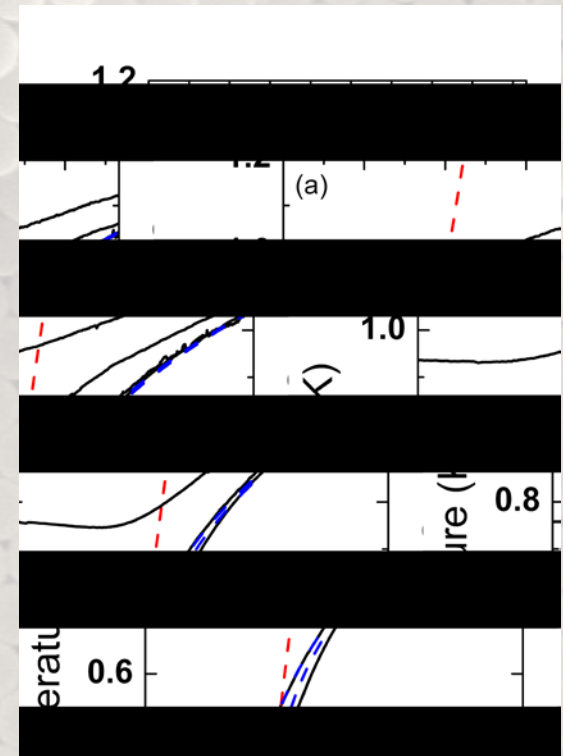
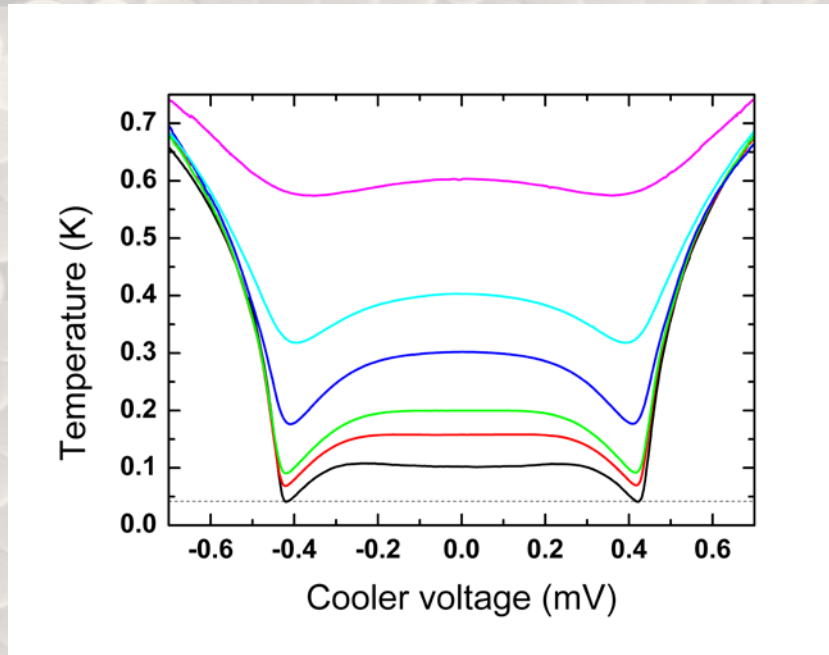
Extremely low G allows measurements of power ~ 10 aW resolution with SINIS thermometry !

[1] C.M. Chang, M.R. Geller, *Phys. Rev. B* **71**, 125304 (2005)

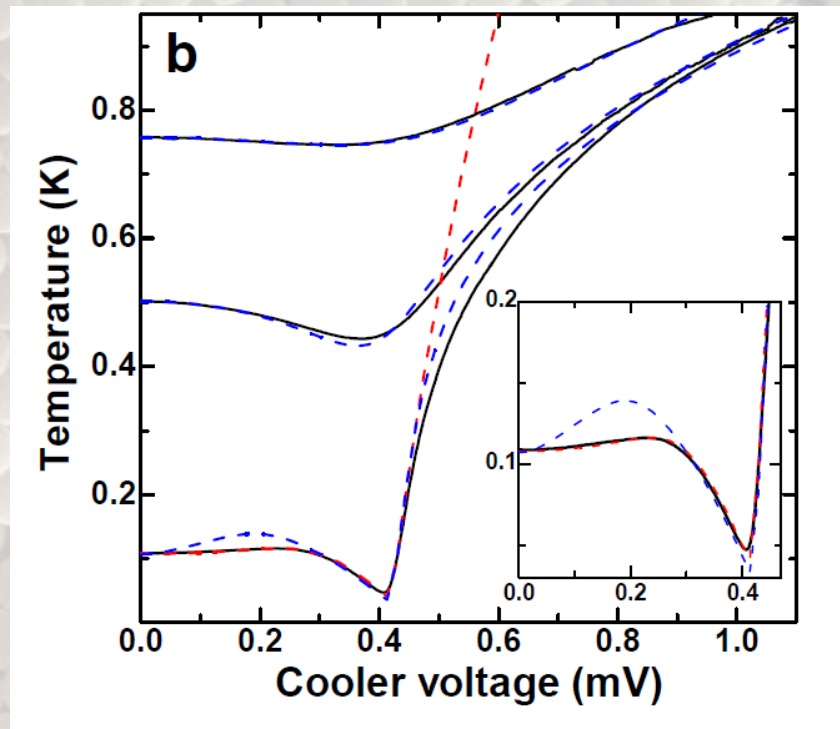
When refrigerator at 50 mK, nanowire
 At 100 mK (noise heating), no cooling
At 42 mK with optimal cooling ($< T_{\text{bath}}$)

For phonon heat flow $P \sim T^n$ ($G \sim T^{n-1}$),
 Data agrees with transition from $n=2.8$
 (red) to $n=6$ (blue)

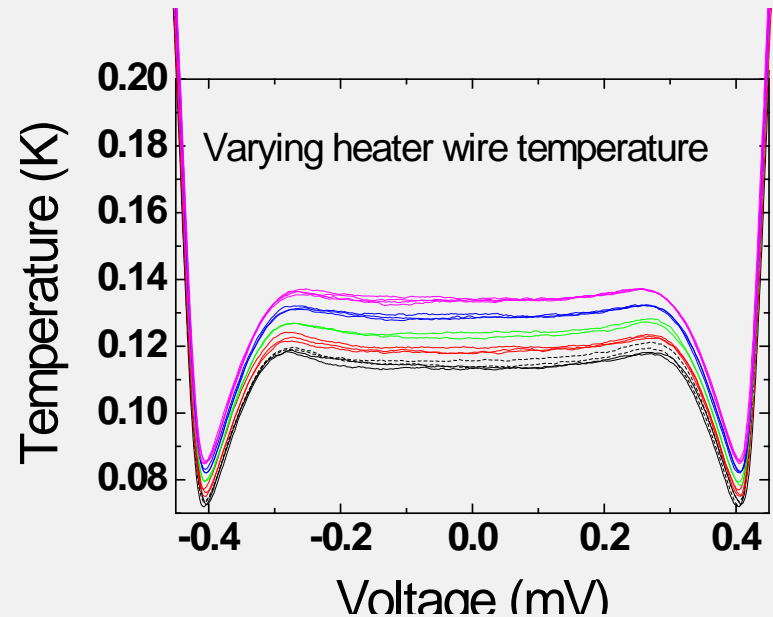
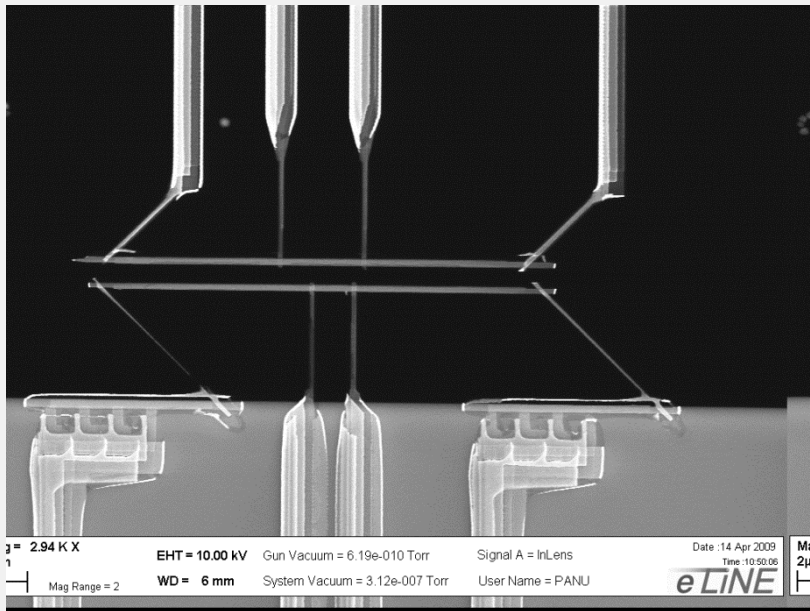
limited by superconductor material (Al)



The point: we can accurately fit the T vs V curve with a thermal model that includes **external input power as a parameter.**

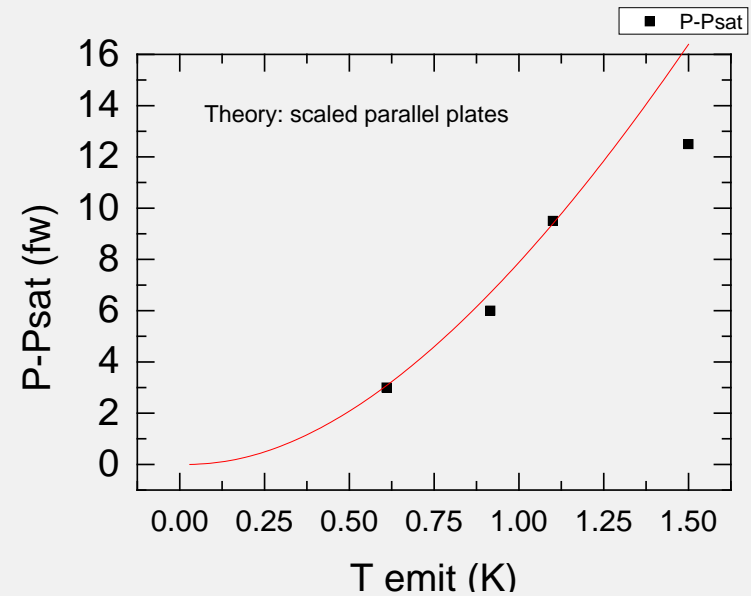
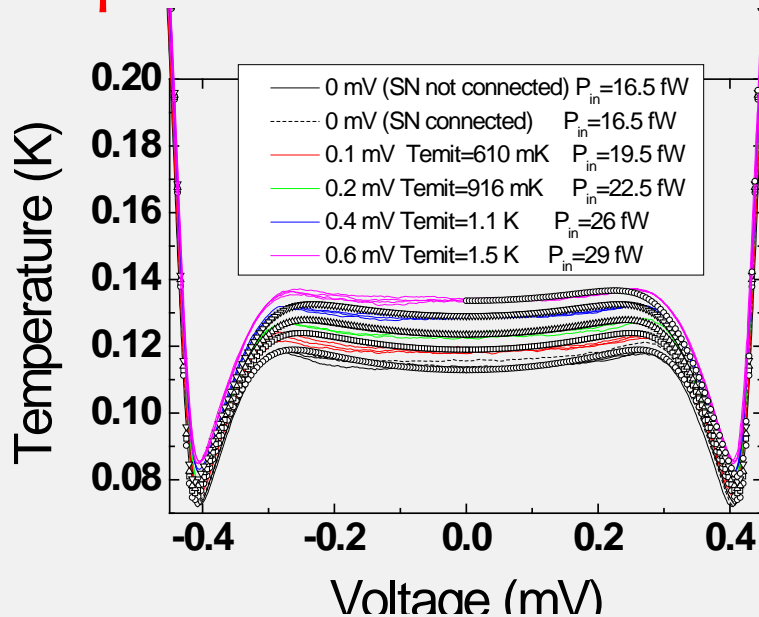


Preliminary data with a suspended cooler as a receiver (gap $\sim 1 \mu\text{m}$)



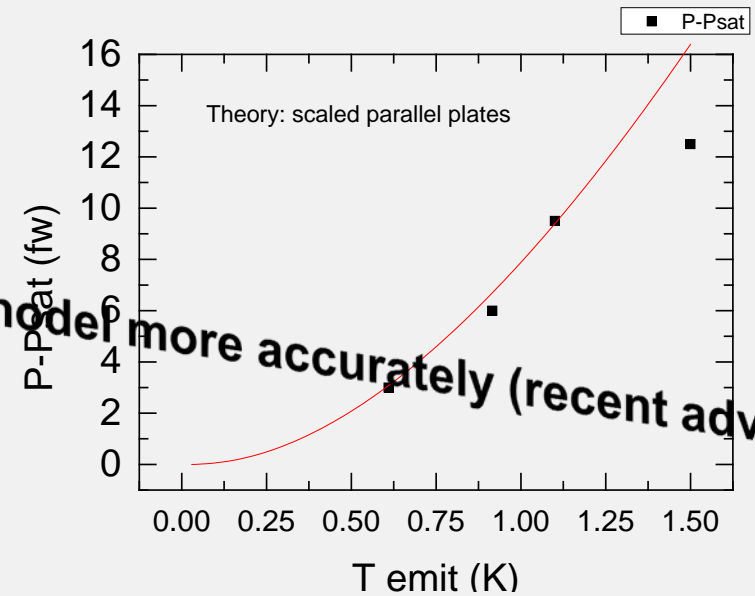
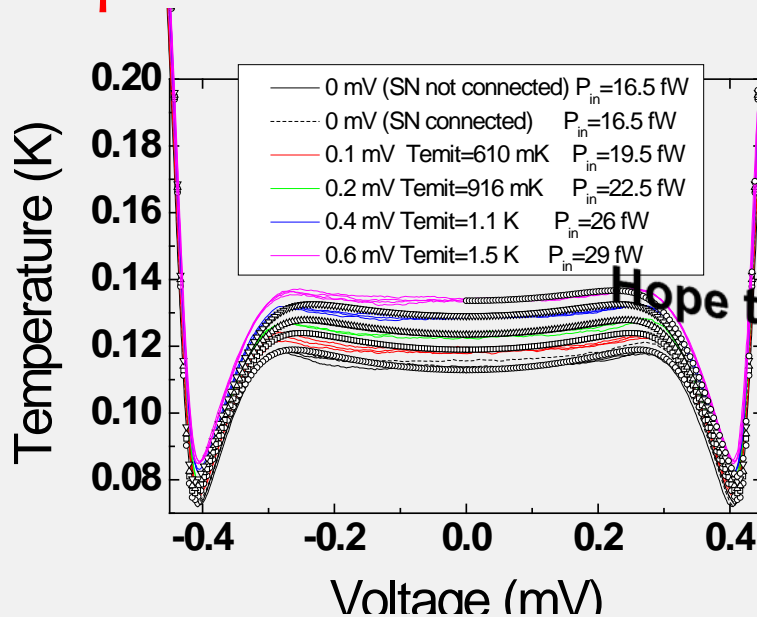
Manuscript in preparation

The point: we can accurately fit the T vs V curve with a thermal model that includes **external input power as a parameter.**



Ten times more power than what parallel plate model seems to predict, but approximately correct temperature dependence ?

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Ten times more power than what parallel plate model seems to predict, but approximately correct temperature dependence ?

Hope to model more accurately (recent advances)

Conclusions:

- Near field thermal radiation can be in a relevant power scale compared with the most sensitive bolometric LT detectors
- At cryogenic temperatures, near field thermal radiation observable up to distances $\sim 10 \mu\text{m}$, currently
- More detailed experiments are in progress and will be finished in the near future

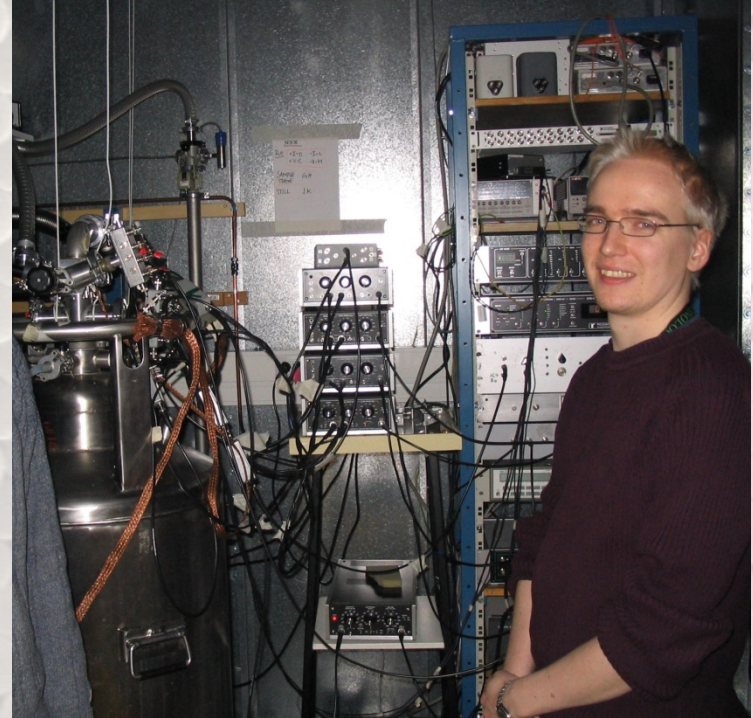
Current group members: Previous group member:



**Dr. Saumyadip
Chaudhuri**



**Tero
Isotalo**



**Dr. Panu Koppinen
(now at VTT Micronova, Espoo)**