

# Fluctuation Electrodynamics – foundations

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thanks to:

Martin Wilkens, Humboldt foundation, DFG



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[www.quantum.physik.uni-potsdam.de](http://www.quantum.physik.uni-potsdam.de)

# Fluctuation electrodynamics

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**what?**

Basic concepts

**how?**

Simple examples

**why not?**

Beyond typical approximations

**got it?**

*Suggested problems*



en.wikipedia.org

(1831–79)



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(1872–1946)



ru.wikipedia.org

Sergei Michailovich Rytov  
(1908–96)

# Fluctuation electrodynamics

– Maxwell equations  $\sim 1861$

$$dF = 0 \quad d^*F = j$$

– Langevin equations  $\sim 1905/08$

$$m \frac{dv}{dt} + \Gamma v = F(t)$$

Langevin force  $F(t) \rightarrow \begin{cases} \text{diffusion} \\ \text{thermal equilibrium} \end{cases}$

Rytov equations  $\nabla \times H - \partial_t D = j(x, t) \leq 1953$

Rytov-Langevin current  $j(x, t) \rightarrow \text{thermal equilibrium}$

- local equilibrium  $T(x)$

Nano-scale

radiative

heat transfer

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atoms  $\ll$  nano  $\ll$   $\lambda$

300 K :  $\lambda \sim 10 \dots 50 \mu\text{m}$

matter  $\sim$  macroscopic continuum

fluctuat'n e'dyn

OK

thermodynamics,  $T$ , entropy  $S$ ,

heat, (ir)reversible processes

non-equilibrium:

driven by external forces  $\nabla T$ ,  $E_L$ ,  $v_2 - v_1$

transfer of momentum = force, torque

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## This lecture

- splash into Maxwell–Langevin equations
- examples: damped oscillator  
hot nano-particle  
heat transfer
- “quantum” remarks

→ *suggested problems*

# Maxwell & Langevin

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macroscopic electrodynamics

$$\partial_t \mathbf{B} = -\nabla \times \mathbf{E} \qquad \nabla \cdot \mathbf{B} = 0$$

$$\partial_t \mathbf{D} = \nabla \times \mathbf{H} - \mathbf{j} \qquad \nabla \cdot \mathbf{D} = \rho$$

linear response of (bulk) material

$$\mathbf{D} = \varepsilon_0 \varepsilon(\mathbf{x}, \omega) \mathbf{E}$$

$$\mathbf{H} = \mu_0^{-1} \mu^{-1}(\mathbf{x}, \omega) \mathbf{B}$$

... always matter that provides nonlinearity

$\varepsilon(\omega), \mu(\omega)$  must be complex (!)

$\varepsilon(\mathbf{x}), \mu(\mathbf{x})$  cannot be local (!)

Rytov & Langevin: losses come with fluctuating sources ('forces')

$$\mathbf{j} = \mathbf{j}_{\text{free}} + \partial_t \mathbf{P}(\mathbf{x}, t) + \nabla \times \mathbf{M}(\mathbf{x}, t)$$

$$\rho = \rho_{\text{free}} - \nabla \cdot \mathbf{P}(\mathbf{x}, t)$$

cf.  $m \frac{dv}{dt} = -\nabla U - \Gamma v + F(t)$

# Maxwell & Langevin

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macroscopic electrodynamics

$$-i\frac{\omega}{c^2}\varepsilon(\mathbf{x},\omega)\mathbf{E} = \nabla \times \mu^{-1}(\mathbf{x},\omega)\mathbf{B} - \mu_0\mathbf{j}(\mathbf{x},\omega) \quad -i\omega\mathbf{B} = \nabla \times \mathbf{E}$$

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Maxwell-Langevin equation: 'stochastic differential equation'

= physicists' code

$$0 = \langle \mathbf{P}(\mathbf{x},t) \rangle$$

word for averages

$$0 \neq \langle \mathbf{P}(\mathbf{x},t)\mathbf{P}(\mathbf{x}',t') \rangle = \int \frac{d\omega}{2\pi} S_P(\mathbf{x},\mathbf{x}',\omega) e^{i\omega(t-t')} \quad \text{spectral density}$$



# Maxwell & Langevin

macroscopic electrodynamics

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Rytov:

$$S_{P,ij}(\mathbf{x},\mathbf{x}',\omega) \approx 2\hbar\bar{N}(\omega)\varepsilon_0 \text{Im} \varepsilon_{ij}(\mathbf{x},\omega)\delta(\mathbf{x} - \mathbf{x}')$$

... fluctuation-dissipation relation

$$\text{Bose-Einstein distribution } \bar{N}(\omega) = \frac{1}{\exp(\hbar\omega/k_B T) - 1} \approx \frac{k_B T}{\hbar\omega}$$

# Maxwell & Langevin

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## *Suggested problems*

### (1) Energy balance in lossy medium

water IR absorption  $\alpha \sim 10^3 \text{ cm}^{-1}$

field  $E = 1 \text{ mV}/\mu\text{m}$ : energy density  $\sim 0.03 \text{ meV}/\mu\text{m}^3$

absorbed power density  $\sim 5 \text{ meV}/\text{ps}/\mu\text{m}^3$

time scale for losses  $c\alpha \sim \omega/2\pi$  frequency

### (2) Temperature: Bjerrum length $\frac{e^2}{\epsilon_0 k_B T} = 0.7 \mu\text{m}$

### (3) Fluctuating dipole moment in volume $\Delta V$ : $d(t) = \int_{\Delta V} dx P(x, t)$

estimate polarization energy  $\frac{\Delta V}{\epsilon_0} \frac{\langle d^2 \rangle_\omega \Delta\omega}{(\Delta V)^2}$

### (4) Count degrees of freedom in volume $1 \mu\text{m}^3$ :

photon modes & matter (phonon) modes (bandwidth  $\Delta\omega$ )

# Maxwell & Langevin

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## Remarks

- equilibrium: fluctuation–dissipation (FD) relation for (quantum) fields Callen & Welton 1951
- FD relations (theorem)
  - Einstein (1905), Langevin (1908): diffusion  $D \leftrightarrow \Gamma$  friction
  - detailed balance: up/down rate  $\sim e^{-\Delta E/k_B T}$  Boltzmann
  - Kubo-Martin-Schwinger (KMS) relation  $\langle A(t)B(t') \rangle_T$  vs  $\langle B(t')A(t) \rangle_T$

# Maxwell & Langevin

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## Remarks

- equilibrium: fluctuation–dissipation (FD) relation for (quantum) fields Callen & Welton 1951
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Einstein (1905), Langevin (1908): diffusion  $D = k_B T \Gamma$  friction  
detailed balance: up/down rate  $\sim e^{-\Delta E/k_B T}$  Boltzmann  
Kubo-Martin-Schwinger (KMS) relation  $S_{AB}(\omega) = e^{-\hbar\omega/k_B T} S_{BA}(\omega)$
- Rytov in full power: non-equilibrium  
*local* temperature  $S_P(\mathbf{x}, \omega) \sim \bar{N}(\omega, T(\mathbf{x})) \text{Im} \varepsilon(\mathbf{x}, \omega)$  Lifshitz & Pitaevskii 1980s
- local equilibrium: on sub- $\lambda$  scale, matter thermalizes ‘fast enough’  
‘photons are a poor thermostat’ Planck 1900 / PLANCK 2013
- local approximation  $\varepsilon(\mathbf{x}, \omega)$  for dielectric response  
non-local media (plasma/screening, effective medium/unit cell)  
genuine surface response  $P(x, y, \omega)\delta(z) \sim \chi(\omega)E(x, y, z = 0, \omega)$

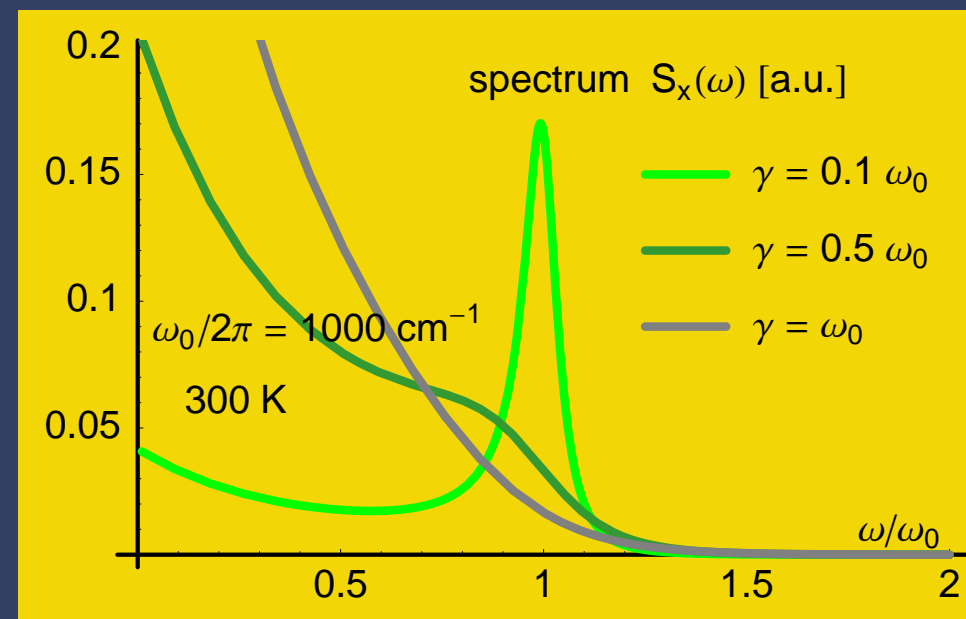
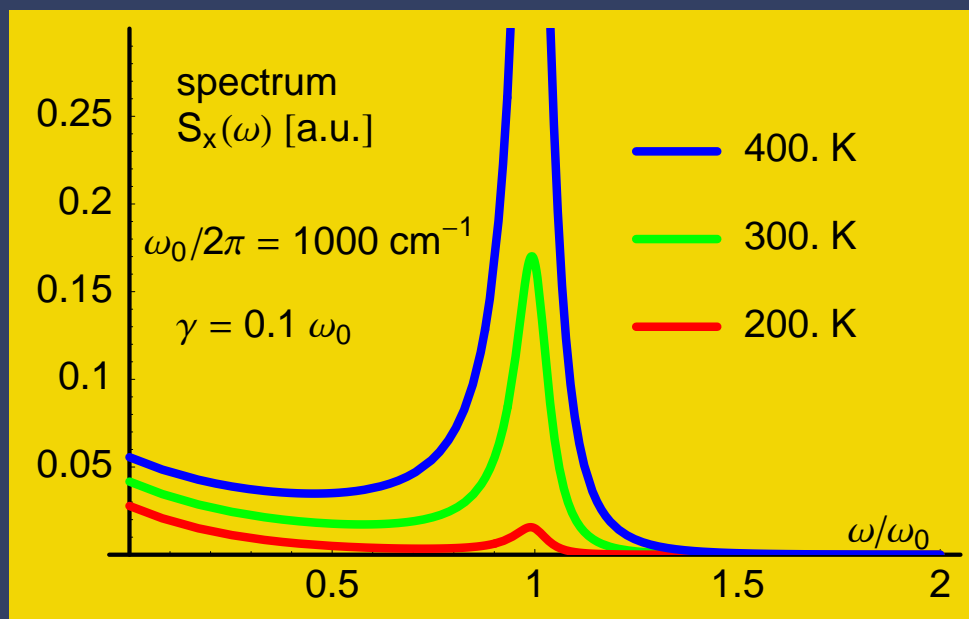
# Examples

*Damped harmonic oscillator*  $m\ddot{x} + \Gamma\dot{x} + Kx = F(t)$

Langevin force spectrum  $S_F(\omega) = 2\hbar\bar{N}(\omega)\omega \operatorname{Re}\Gamma(\omega)$  white vs colored (discussion)

Oscillator spectrum (equilibrium, long-time limit)

$$S_x(\omega) = \frac{S_F(\omega)}{|-m\omega^2 - i\omega\Gamma(\omega) + K|^2} = 2\hbar\bar{N}(\omega) \operatorname{Im}\left(\frac{1}{-m\omega^2 - i\omega\Gamma(\omega) + K}\right)$$



non-equilibrium: relaxation dynamics, two-temperature driving ...

quantum questions: zero-point energy, negative frequencies, overdamped limit ...

# Examples

## Hot nano-sphere



spectrum of dipole moment

$$S_d(\omega) \approx 2\hbar\bar{N}(\omega)4\pi a^3\epsilon_0 \operatorname{Im} \frac{\epsilon(\omega) - 1}{\epsilon(\omega) + 2}$$

polarizability  $\operatorname{Im} \alpha(\omega)$

• emitted field

$$\mathbf{E}(\mathbf{r}, \omega) = \frac{\omega^2 e^{i\omega r/c}}{4\pi\epsilon_0 c^2 r} \left\{ (\mathbf{d} - 3\hat{\mathbf{r}}(\hat{\mathbf{r}} \cdot \mathbf{d})) \left( \frac{-1}{(\omega r/c)^2} + \frac{\mathbf{i}}{\omega r/c} \right) + \mathbf{d} - \hat{\mathbf{r}}(\hat{\mathbf{r}} \cdot \mathbf{d}) \right\}$$

$$\langle \mathbf{E}(\mathbf{r}, \omega) \rangle = 0 = \mathbf{G}(\mathbf{r}, 0, \omega) \cdot \mathbf{d}(\omega)$$

Green function/tensor/dyadic  
→ this week

# Examples

## Hot nano-sphere



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Green function/tensor/dyadic  
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## Suggested problems

average intensity “ $\langle |\mathbf{E}(\mathbf{r}, \omega)|^2 \rangle$ ” vs distance  $r$

average Poynting vector “ $\operatorname{Re} \langle \mathbf{E}^*(\mathbf{r}, \omega) \times \mathbf{H}(\mathbf{r}, \omega) \rangle$ ”

total emitted power & cooling rate/time

array of nano-particles

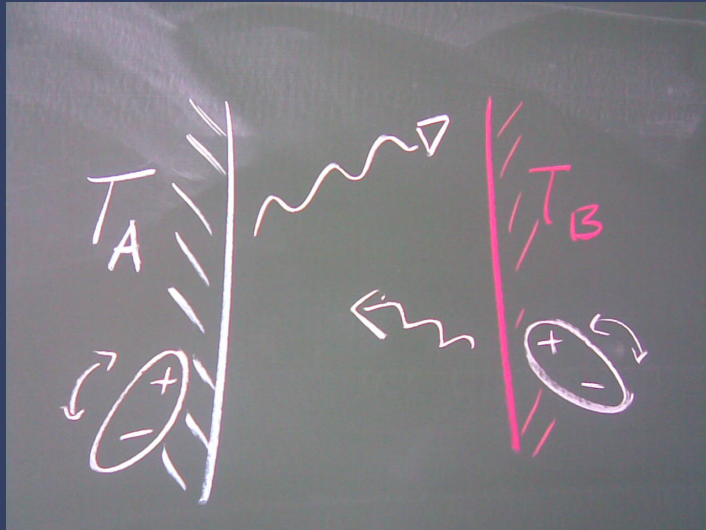
→ P Ben-Abdallah

- core technique for non-equilibrium Rytov theory  
= ‘incoherent summation’ over elementary source volumes

# Examples

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## Radiative heat transfer



Poynting vector  $\text{Re} \langle \mathbf{E}^*(\mathbf{r}, \omega) \times \mathbf{H}(\mathbf{r}, \omega) \rangle_z$

$T_{z0}$  current

= (left sources) + (right sources)

$\approx (T_B - T_A) \times h(\text{distance}, \bar{T})$

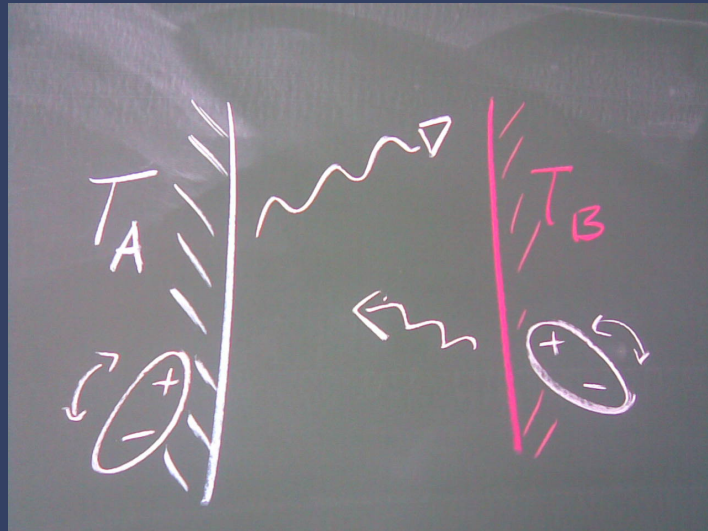
→ this week's programme

- Casimir interaction = momentum transfer, stress tensor  $\langle T_{ij} \rangle$



# Examples

## Radiative heat transfer



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## Typical / tacit assumptions

- homogeneous temperature profile *suggested problem* convection faster than radiation
- Langevin sources not correlated radiative coupling slower than local relaxation
- no heat extracted from vacuum fluctuations but: unstable vacua

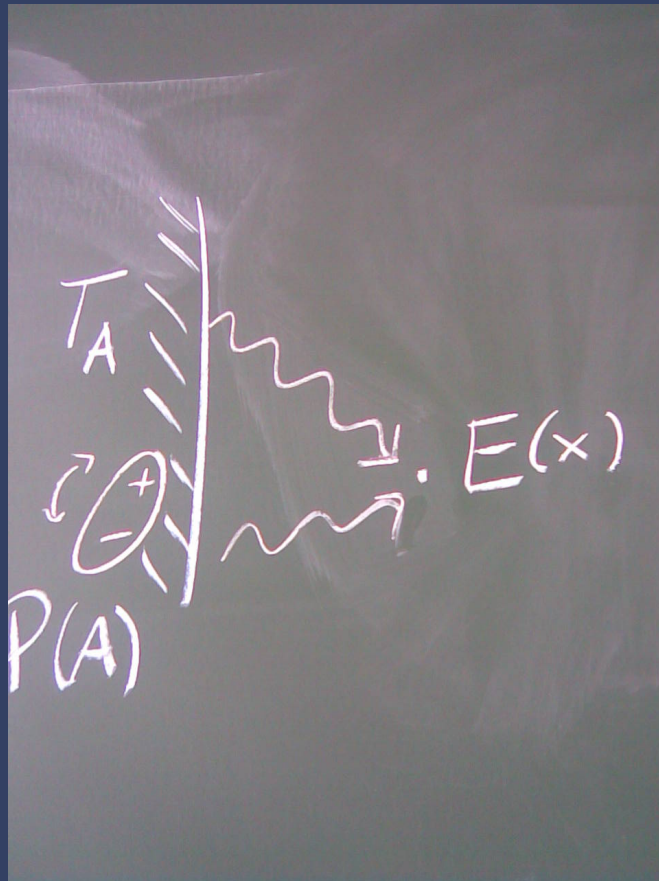
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# Examples

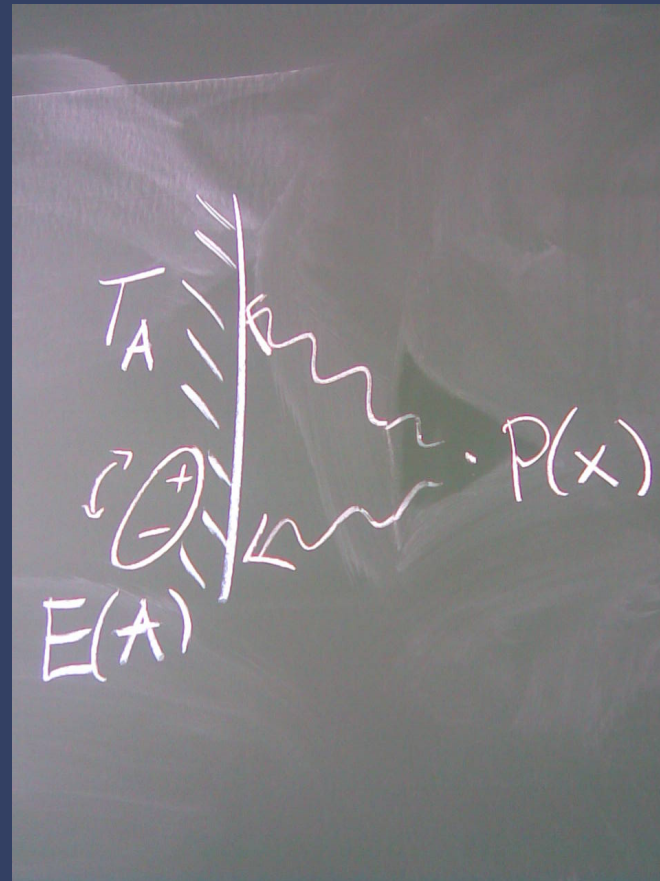
## Radiative heat transfer

- convenient shortcut: generalized Kirchhoff law

Dorofeyev & Vinogradov (*Phys Rep* 2011)

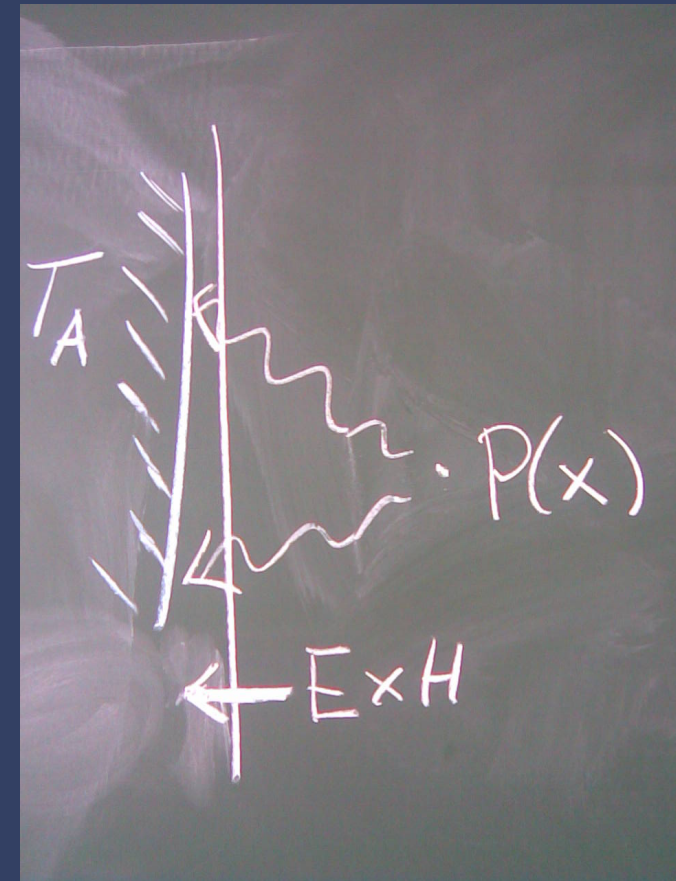


emission by sources



reciprocity

$$G(x, x') = G^T(x', x)$$



Poynting vector

# Examples

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# “Quantum” remarks

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(1) ‘first quantization’    QM            electrodynamics

Heisenberg                     $\Delta p \Delta x \sim \hbar$     (Fourier)

Planck & Einstein             $\hbar\omega$                      $\epsilon E^2 + \mu H^2$

de Broglie                     $\hbar k$                      $\epsilon E \times B$

(2) ‘second quantization’

Pauli & Jordan                     $\Delta E \Delta A \sim \hbar$

$$i[A_i(\mathbf{x}), E_j(\mathbf{x}')] = \frac{\hbar}{\epsilon_0} \delta_{ij}^\perp(\mathbf{x} - \mathbf{x}')$$

$$i[B_i(\mathbf{x}), E_j(\mathbf{x}')] = -\frac{\hbar}{\epsilon_0} \epsilon_{ijk} \partial_k \delta(\mathbf{x} - \mathbf{x}')$$

## Ordering of operators

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Feynman             $\langle T\{E(t)E(t')\} \rangle$              $D_E(\omega)$  even

Green                 $i\langle [E(t), E(t')] \rangle \Theta(t - t')$      $\text{Im } G(\omega)$  odd

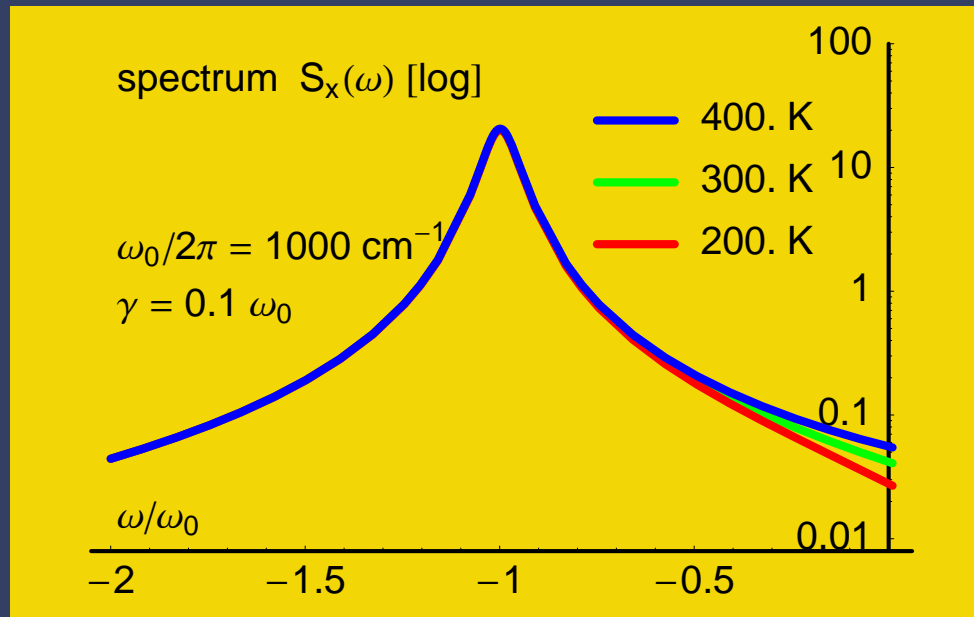
Hadamard             $\langle \{E(t), E(t')\} \rangle$              $H_E(\omega)$  real and even     $\sim \bar{N}(\omega) + \frac{1}{2}$

Glauber (?)             $\langle E(t)E(t') \rangle$              $S_E(\omega)$  positive             $\sim \bar{N}(\omega), \quad \omega > 0$   
 $e^{i\omega(t-t')}$

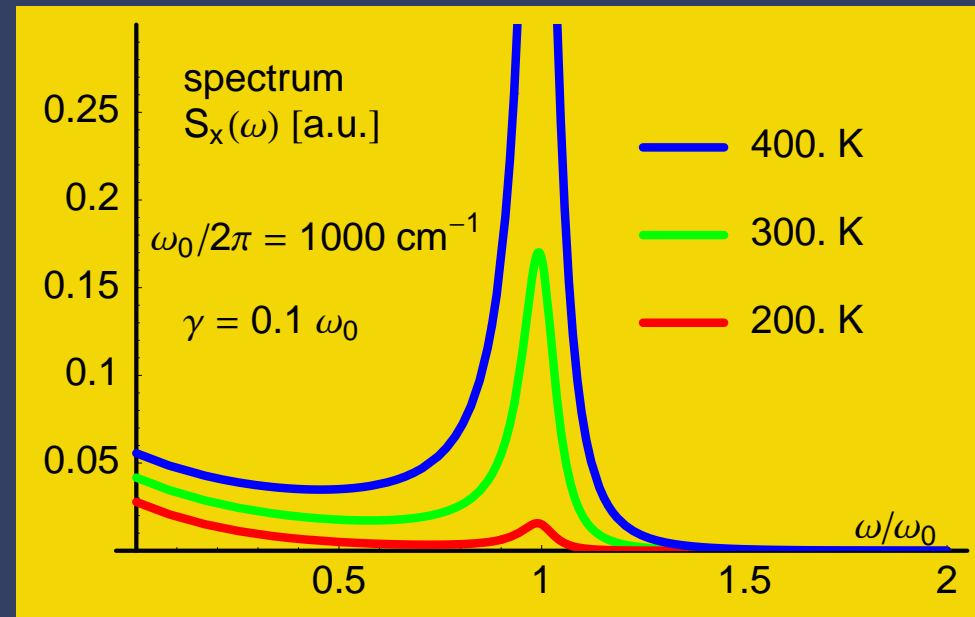
# "Quantum" remarks

## Damped oscillator in equilibrium

Glauber  $S_x(\omega)$  at negative ...



... and positive frequencies

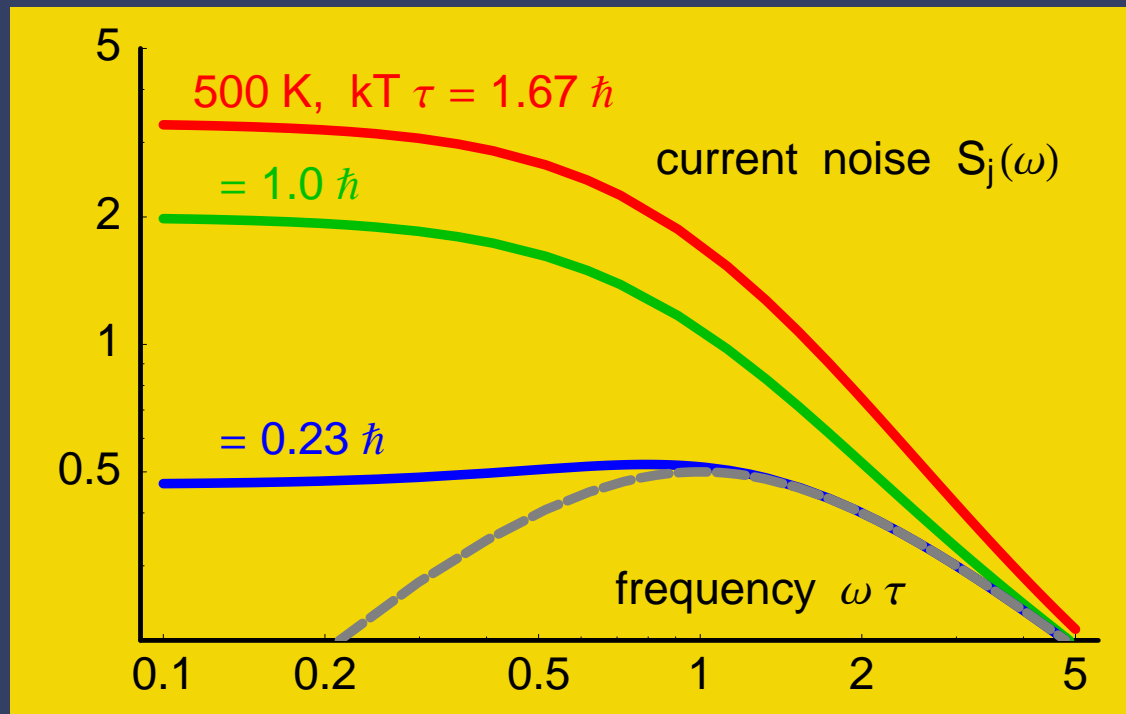


quantum noise / zero-point fluctuations

# “Quantum” remarks

## Electric current noise in metal

$$\text{symmetrized } H_j(\omega) \sim (\bar{N}(\omega) + \frac{1}{2}) \omega \text{Re} \sigma(\omega)$$



quantum noise of overdamped/diffusive field

- Casimir interaction: see Intravaia & CH (*Phys Rev Lett* 2009)

# “Quantum” remarks

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## *Suggested problems*

Glauber spectrum:  $-\bar{N}(-\omega) = \bar{N}(\omega) + 1$  asymmetric, still positive

Magnetic flux quantization? usually need electron charge for flux quantum, so it's beyond field quantization? (cf. magnetic monopole)

Why do annihilation operators evolve with positive frequencies?

$$a(t) = a(0) \exp(-i\omega t)$$

Can all observers agree on the sign of a frequency? (not always!)

anomalous Doppler effect (Ginzburg):

$$\omega' = \gamma(\omega + \mathbf{v} \cdot \mathbf{k}) < 0$$

quantum friction, pair creation, Cherenkov radiation

# Summary

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**stochastic electrodynamics**    Rytov = Maxwell + Langevin

route towards thermal equilibrium

QED: preserve commutators

separation of time scales

field absorption vs heat transport

## Examples

stochastic field correlations: linear in sources

natural non-equilibrium technique (local vs global eq)

'macroscopic' response functions sufficient, even for quantum noise

( $\epsilon$ ,  $\mu$ , surface response ...)

## Glimpses ... of the land beyond Rytov

*Suggested problem*

wikipedia entries about S. M. Rytov

(english, français, deutsch ...)



# Thanks to the group!

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Francesco Intravaia (→ Nottingham)	Casimir physics	<i>Phys Rev A</i> 2010
Jürgen Schiefele (→ Madrid)	atom chip & BEC	<i>Phys Lett A</i> 2011
Harald Haakh (→ Erlangen)	superconducting atom chips	<i>Eur Phys J B</i> 2012
Geesche Boedecker (PhD)	non-Markovian QED	<i>Ann Phys (Berlin)</i> 2013
Giuseppe Cammarata (post doc)	quantization of coupled oscillators	
Alexander Kegeles (PhD)	entanglement production	
Gregor Pieplow (Dipl)	quantum friction	<i>New J Phys</i> 2013
Ralf Saplata (Dipl)	failures of adiabaticity	
Abdoulaye Diallo (Dipl)	border of BEC	
Timo Felbinger	Illarion Dorofeyev (Nizhny Novgorod)	
Martin Wilkens	Vanik E. Mkrtchian (Ashtarak)	

## Appendix – forgotten references

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## Appendix — Fluctuation-dissipation relation

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Cross correlation spectrum:

$$\mathcal{E}_{ij}(\mathbf{x}, \mathbf{x}'; \omega) = \int_{-\infty}^{+\infty} d\tau e^{-i\omega\tau} \langle E_i(\mathbf{x}, t + \tau) E_j(\mathbf{x}', t) \rangle$$

Linear response:

$$\mathcal{G}_{ij}(\mathbf{x}, \mathbf{x}'; \omega) = \frac{i}{\hbar} \int_0^{+\infty} d\tau e^{i\omega\tau} \langle [E_i(\mathbf{x}, \tau), E_j(\mathbf{x}', 0)] \rangle$$

Canonical ensemble:  
(KMS relation)

$$\int_{-\infty}^{+\infty} d\tau e^{i\omega\tau} \langle A(\tau) B(0) \rangle = e^{\hbar\omega/k_B T} \int_{-\infty}^{+\infty} d\tau e^{i\omega\tau} \langle B(0) A(\tau) \rangle$$

→ Fluctuation–dissipation relation:

$$\mathcal{E}_{ij}(\mathbf{x}', \mathbf{x}; \omega) = 2\hbar\bar{N}(\omega) \frac{\mathcal{G}_{ji}(\mathbf{x}', \mathbf{x}; \omega) - [\mathcal{G}_{ij}(\mathbf{x}, \mathbf{x}'; \omega)]^*}{2i}$$

Reciprocity:  $\mathcal{G}_{ji}(\mathbf{x}', \mathbf{x}; \omega) = \mathcal{G}_{ij}(\mathbf{x}, \mathbf{x}'; \omega)$